

Investigation of fundamental properties of the $\bar{K}NN$ state

32nd J-PARC PAC meeting
(14-16, July, 2021)

Takumi Yamaga (RIKEN)

Short summary

– For stage-1 approval –

Beam-time request

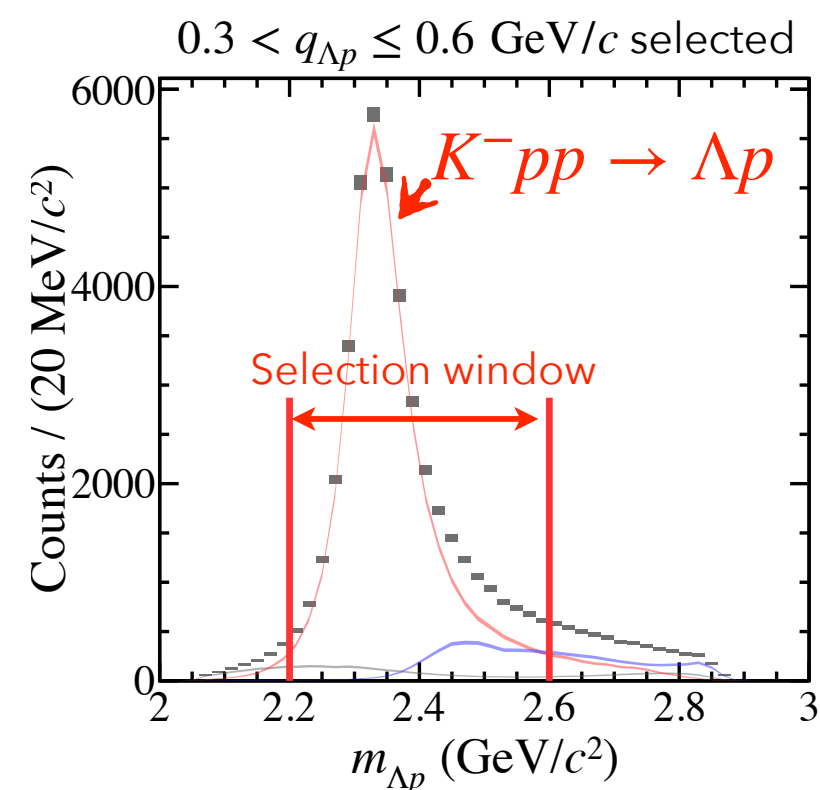
8 weeks $\otimes$ 90 kW @ K1.8BR

(648 kW·week @ 90% up-time)

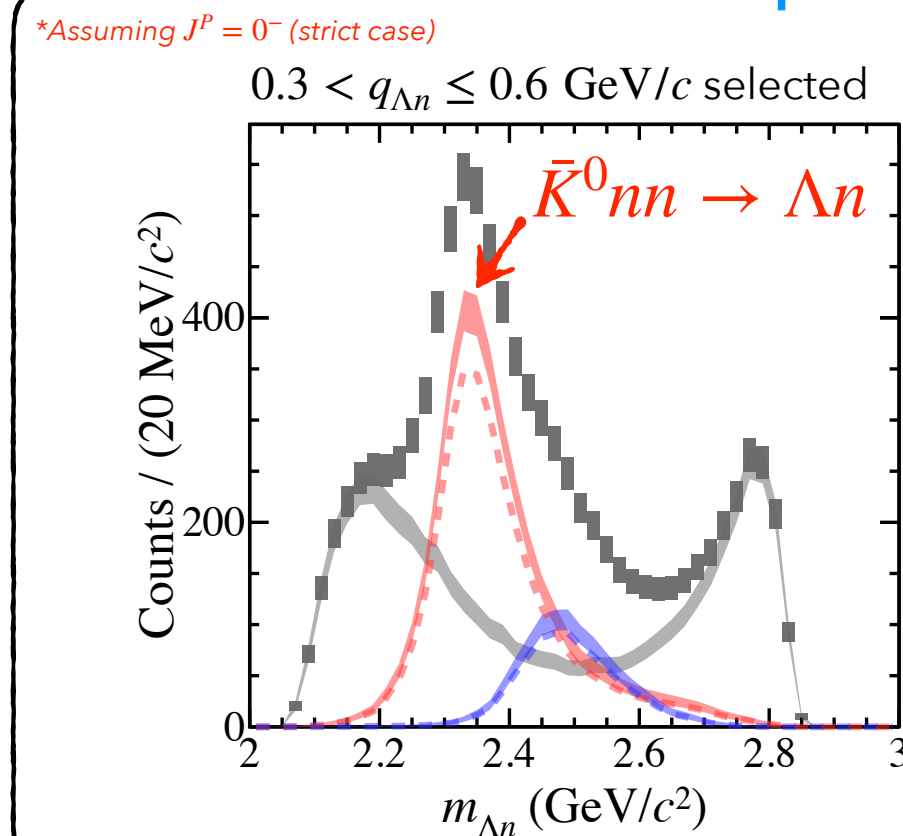
Expected results

To confirm $I = 1/2$ $\bar{K}NN$ doublet

Measurement for Λp pair

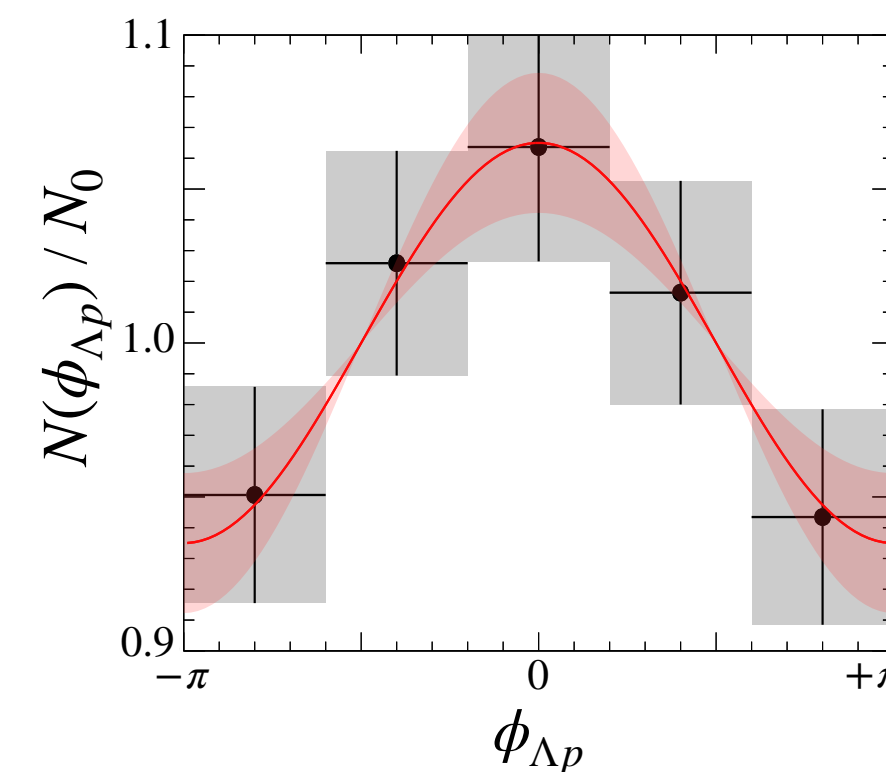


Measurement for Λn pair

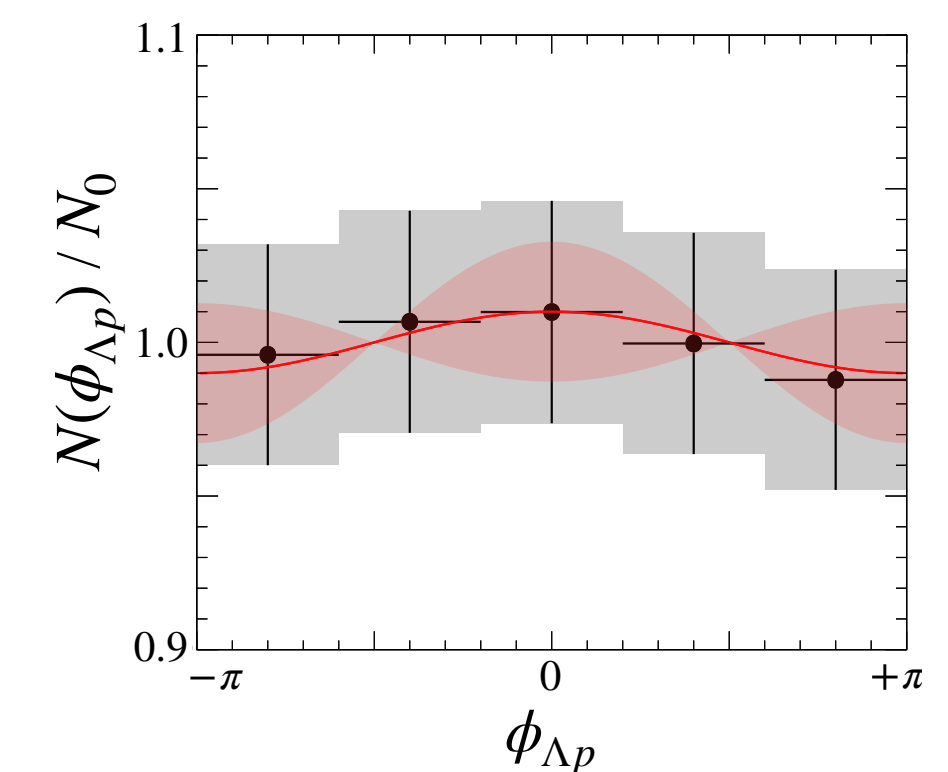


To determine J^P of $\bar{K}NN$

$J^P = 0^-$



$J^P = 1^-$

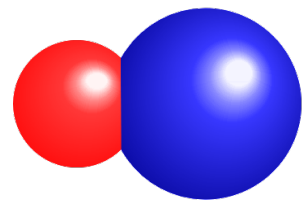


New programs for kaonic nuclei

– Further investigation of $\bar{K}NN$ & Searching for lighter & heavier systems –

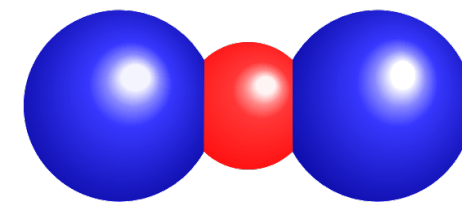
– *This experiment* –

Lighter system



$$\Lambda(1405) = K^- p$$

$\bar{K}NN$ system



Most fundamental / Internal structure

J^P determination

Study for $\bar{K}^0 nn$

Relation to Λ^*

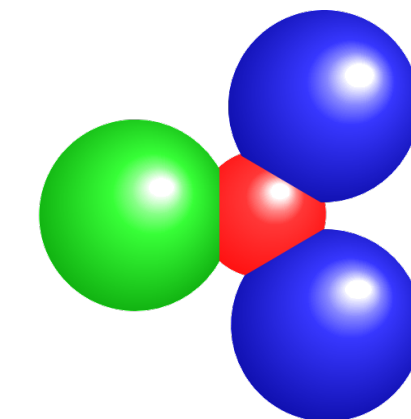
Decay branch

Large Γ

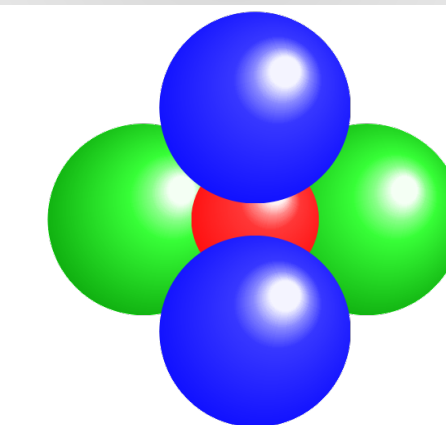
Heavier system

E80 $\bar{K}NNN$ system

Spokesperson : F. Sakuma



$\bar{K}NNNN$ system



$\bar{K}\alpha\alpha$ system

The simplest kaonic nucleus system, $\bar{K}NN$

– Exotic bound state resulting strong attractive $\bar{K}N$ interaction –

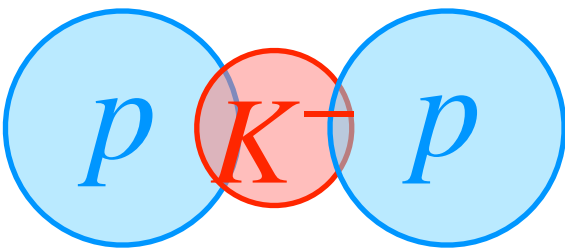
$\bar{K}NN$ bound state

Bound system of
anti-kaon and two nucleons

$$\left[\bar{K}_{I=\frac{1}{2}} (NN)_{I=1} \right]_{I=\frac{1}{2}} \quad (J^P = 0^-)$$

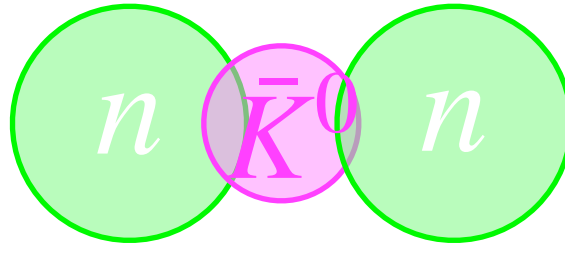
$$I_{\bar{K}NN}^{(z)} = +1/2 \quad \text{“} K^- pp \text{”}$$

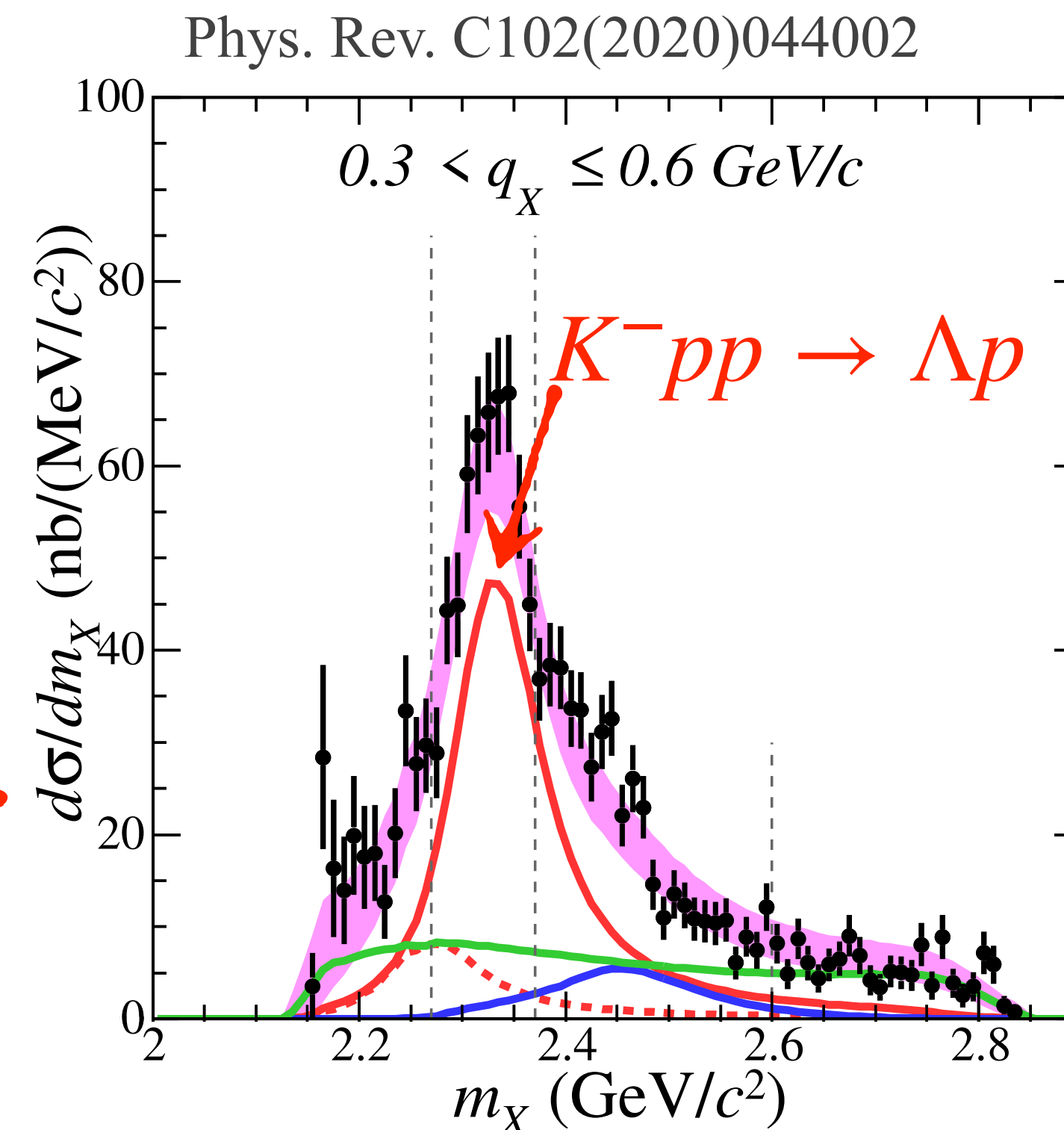
$(K^- pp - \bar{K}^0 pn)$



$$I_{\bar{K}NN}^{(z)} = -1/2 \quad \text{“} \bar{K}^0 nn \text{”}$$

$(\bar{K}^0 nn - K^- pn)$

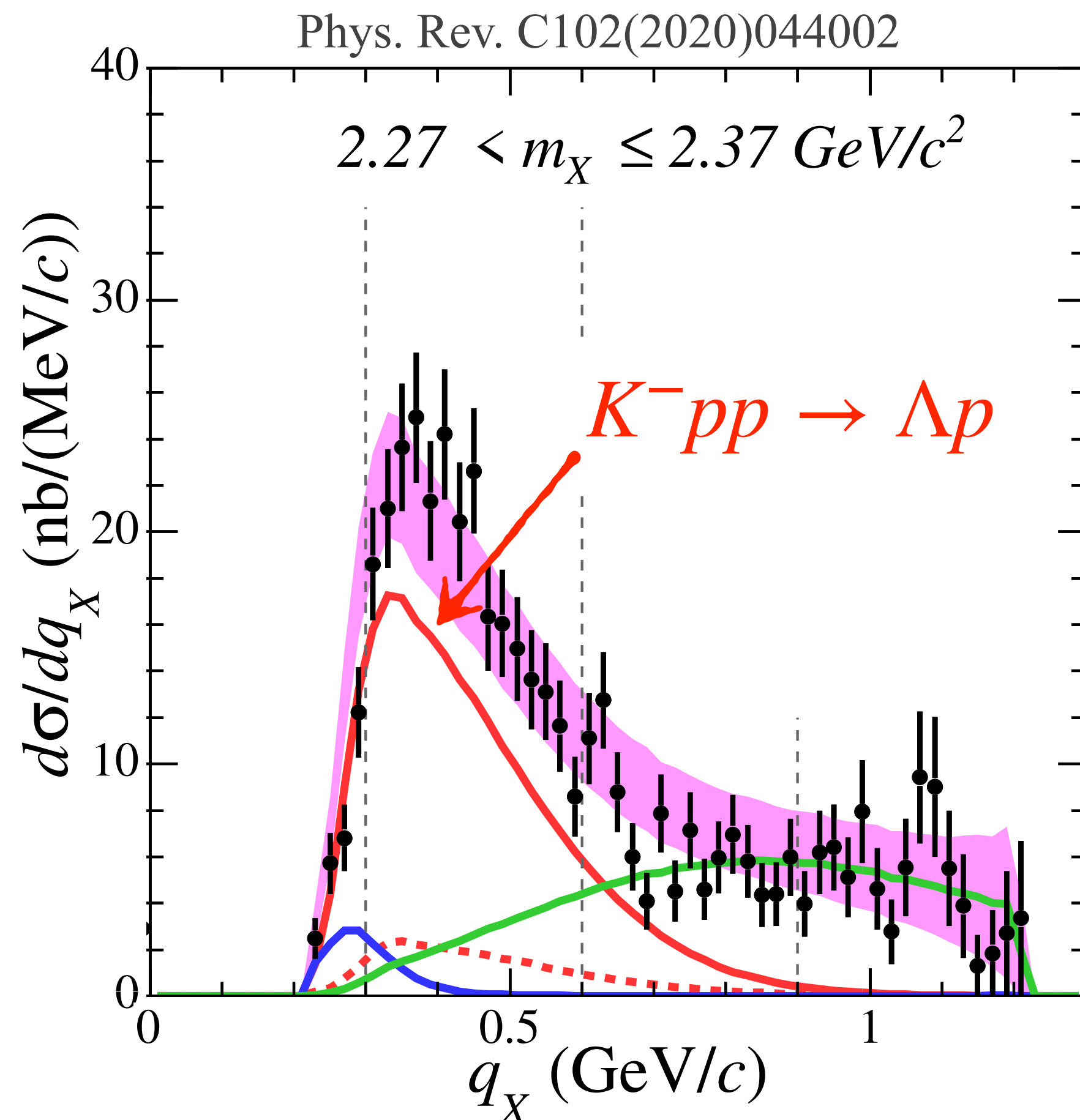




A clear formation signal of $K^- pp$ was observed.

The simplest kaonic nucleus system, $\bar{K}NN$

– Momentum transfer dependence & spacial size of the $\bar{K}NN$ –



q_X -dependence based on PWIA

$$\frac{d\sigma}{dq_X} \propto \exp\left(-\frac{q^2}{Q^2}\right)$$

$Q \sim 400 \text{ MeV}/c$: S-wave Gaussian form factor

The result suggests that the spacial size would be surprisingly small ($r \sim 0.6 \text{ fm}$).

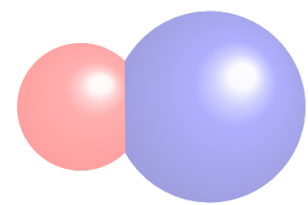
We need more precise measurement & analysis to conclude the spacial size of the $\bar{K}NN$

New programs for kaonic nuclei

– Further investigation of $\bar{K}NN$ & Searching for lighter & heavier systems –

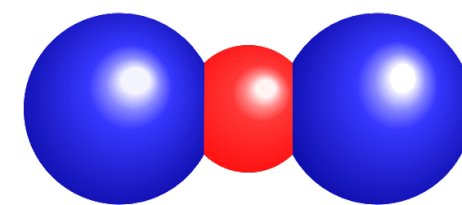
– *This experiment* –

Lighter system



$$\Lambda(1405) = K^- p$$

$\bar{K}NN$ system



Most fundamental / Internal structure

J^P determination

Study for $\bar{K}^0 nn$

Relation to Λ^*

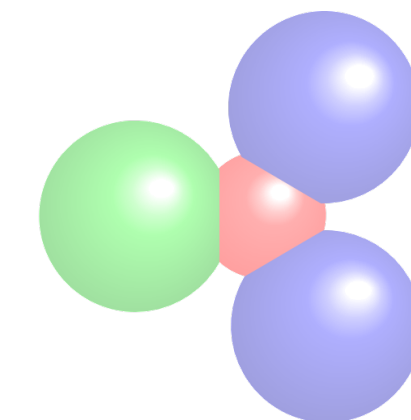
Decay branch

Large Γ

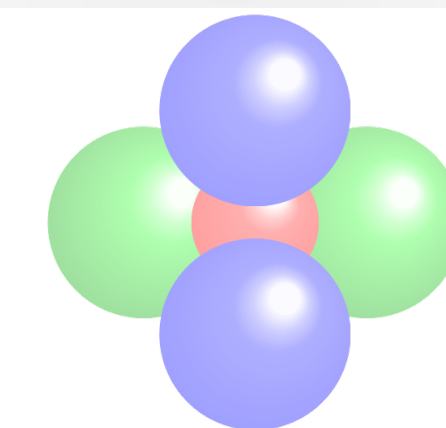
Heavier system

E80 $\bar{K}NNN$ system

Spokesperson : F. Sakuma



$\bar{K}NNNN$ system



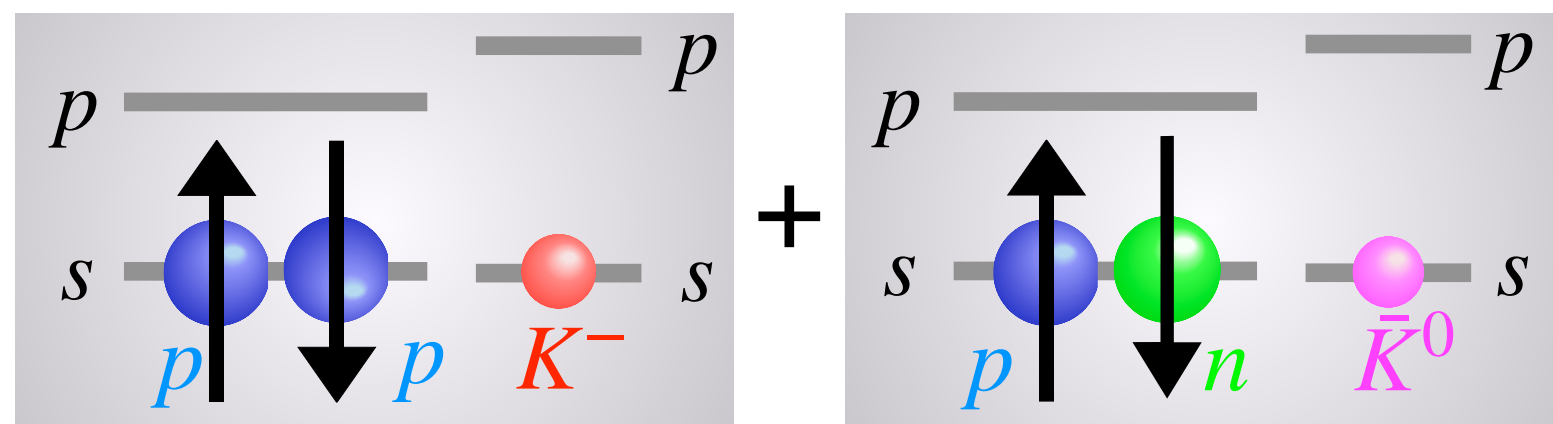
$\bar{K}\alpha\alpha$ system

Internal structure & spin-parity

There are two possible J^P as for the $\bar{K}NN$ ground state.

" $(NN)_{(I.sym \times S.asym)} \otimes \bar{K}$ "

$$J^P = 0^-$$

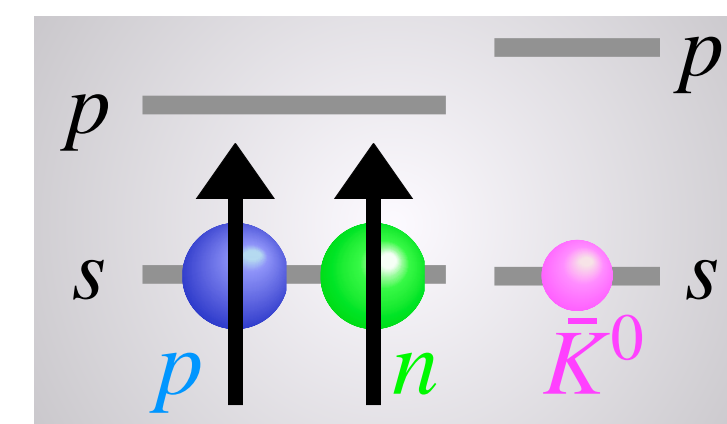


$$\frac{|I_{\bar{K}N} = 0|}{|I_{\bar{K}N} = 1|} = \frac{3}{1}$$

Deeper bound expected

" $(NN)_{(I.asym \times S.sym)} \otimes \bar{K}$ "

$$J^P = 1^-$$



$$\frac{|I_{\bar{K}N} = 0|}{|I_{\bar{K}N} = 1|} = \frac{1}{3}$$

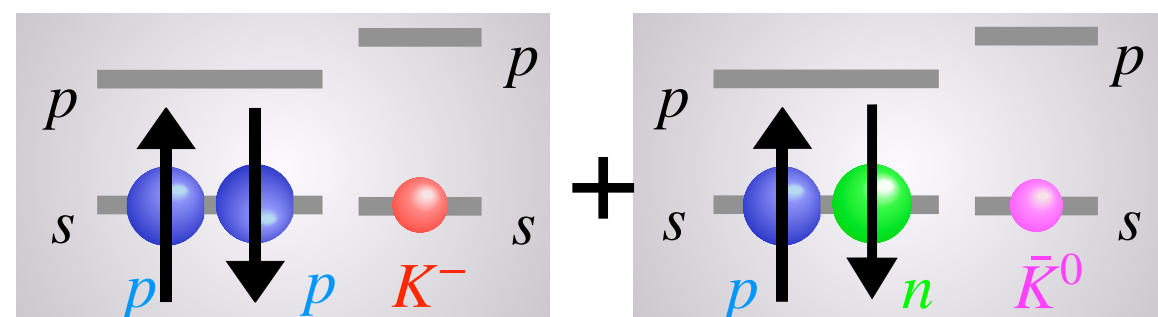
Shallower bound expected

* Positive parity state should be higher excited state if exist.

How to determine J^P

– Λp spin-spin correlation ($\alpha_{\Lambda p}$) in $K^- pp \rightarrow \Lambda p$ decay –

$$J^P = 0^-$$

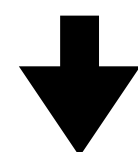


To make negative parity from Λp

$$L_{\Lambda p} = 1$$

To be $J = 0$

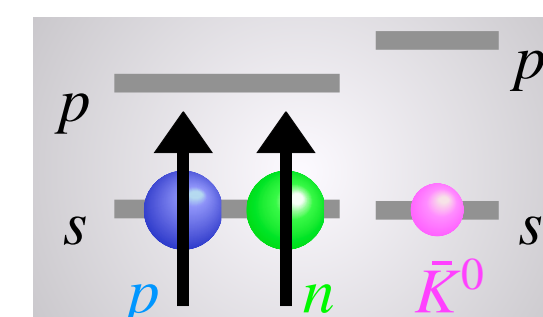
$$S_{\Lambda p} = 1$$



$$\alpha_{\Lambda p} = +1$$

Spin parallel

$$J^P = 1^-$$



To make negative parity from Λp

$$L_{\Lambda p} = 1$$

To be $J = 1$

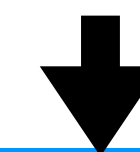
$BR = 1/3$

$$S_{\Lambda p} = 0$$

+

$$S_{\Lambda p} = 1$$

$BR = 2/3$



$$\alpha_{\Lambda p} = -1$$

Spin anti-parallel



$$\alpha_{\Lambda p} = +1$$

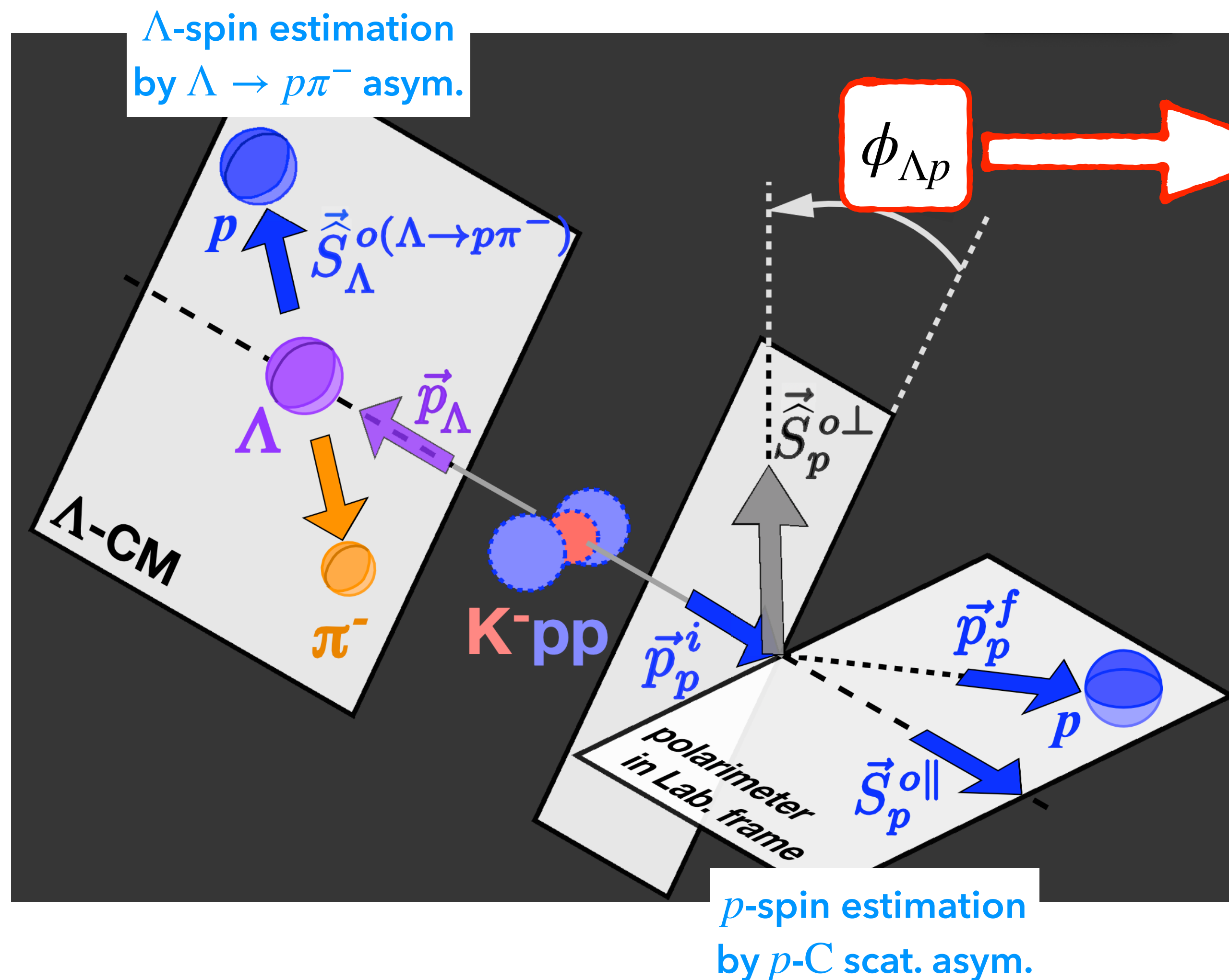
Spin parallel

We can deduce J^P from $\alpha_{\Lambda p}$ measurement.

Experimental method

How to measure spin-spin correlation

– Spin alignment measurement by $\Lambda \rightarrow p\pi^-$ & p -C scattering –



Spin-spin correlation on ϕ -asymmetry

$$N(\phi_{\Lambda p}) = N_0 \cdot (1 + r^{(J^P)} \cdot \alpha_{\Lambda p} \cos \phi_{\Lambda p})$$

$r^{(J^P)}$: asymmetry reduction factor defined by;

α_- : Λ asym. parameter B : Magnetic field

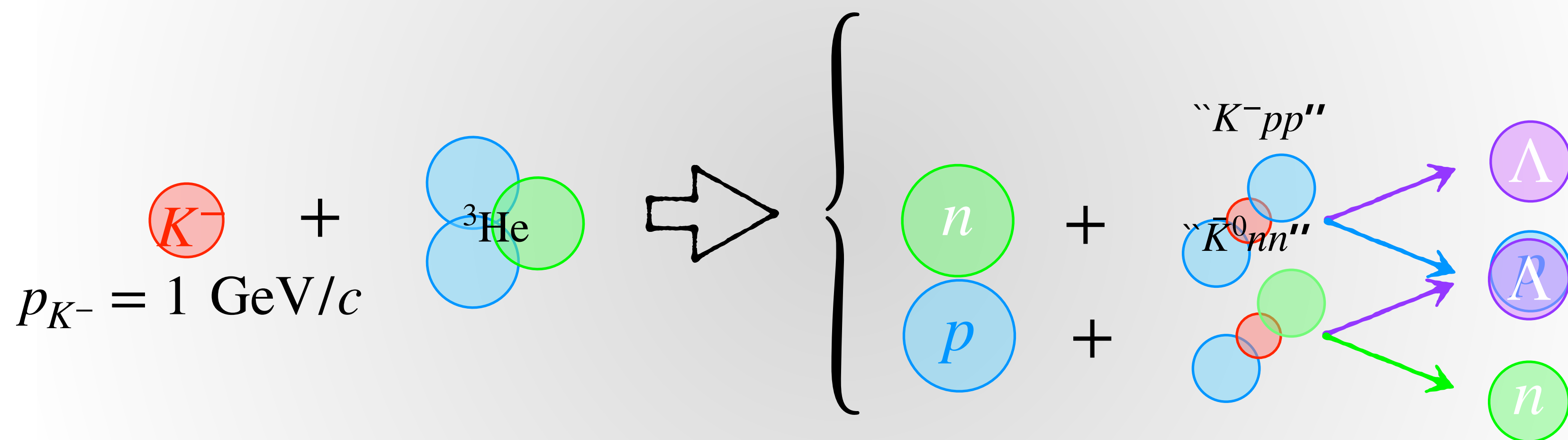
A_{pC} : Analyzing power $B_{\bar{K}}$: Binding energy

$f_{\vec{S}_\Lambda}$: Spin distribution q : Momentum transfer

We can observe $\alpha_{\Lambda p}$ from $\phi_{\Lambda p}$ -distribution.

How to produce $\bar{K}NN$

$\bar{K}NN$ production in $K^- + {}^3\text{He} \rightarrow \Lambda pn$ reaction



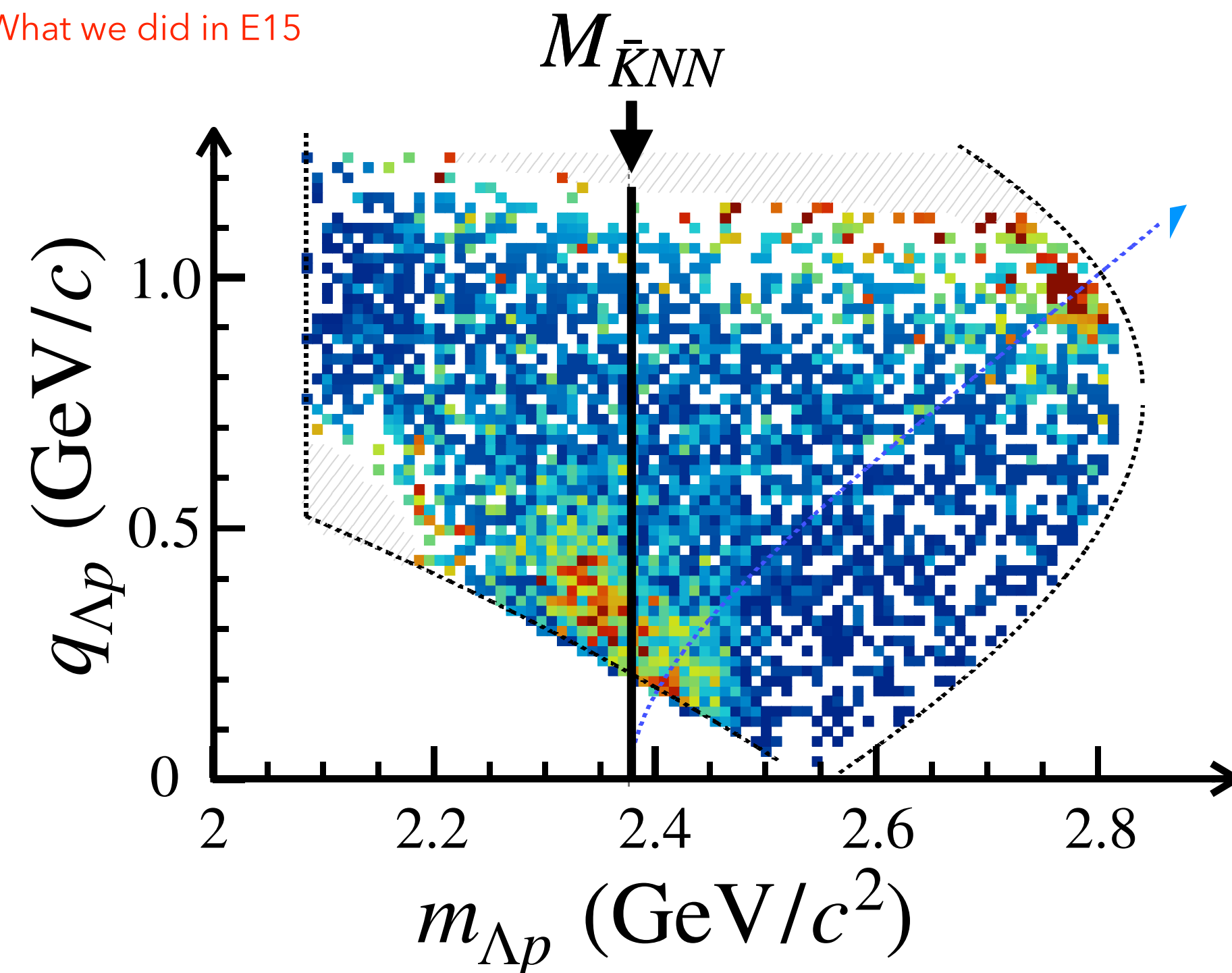
Both K^-pp & $\bar{K}^0nn$ are produced in $K^- + {}^3\text{He}$ reaction.

How to measure $\bar{K}NN$

– 2D analysis for invariant-mass & momentum transfer of ΛN systems –

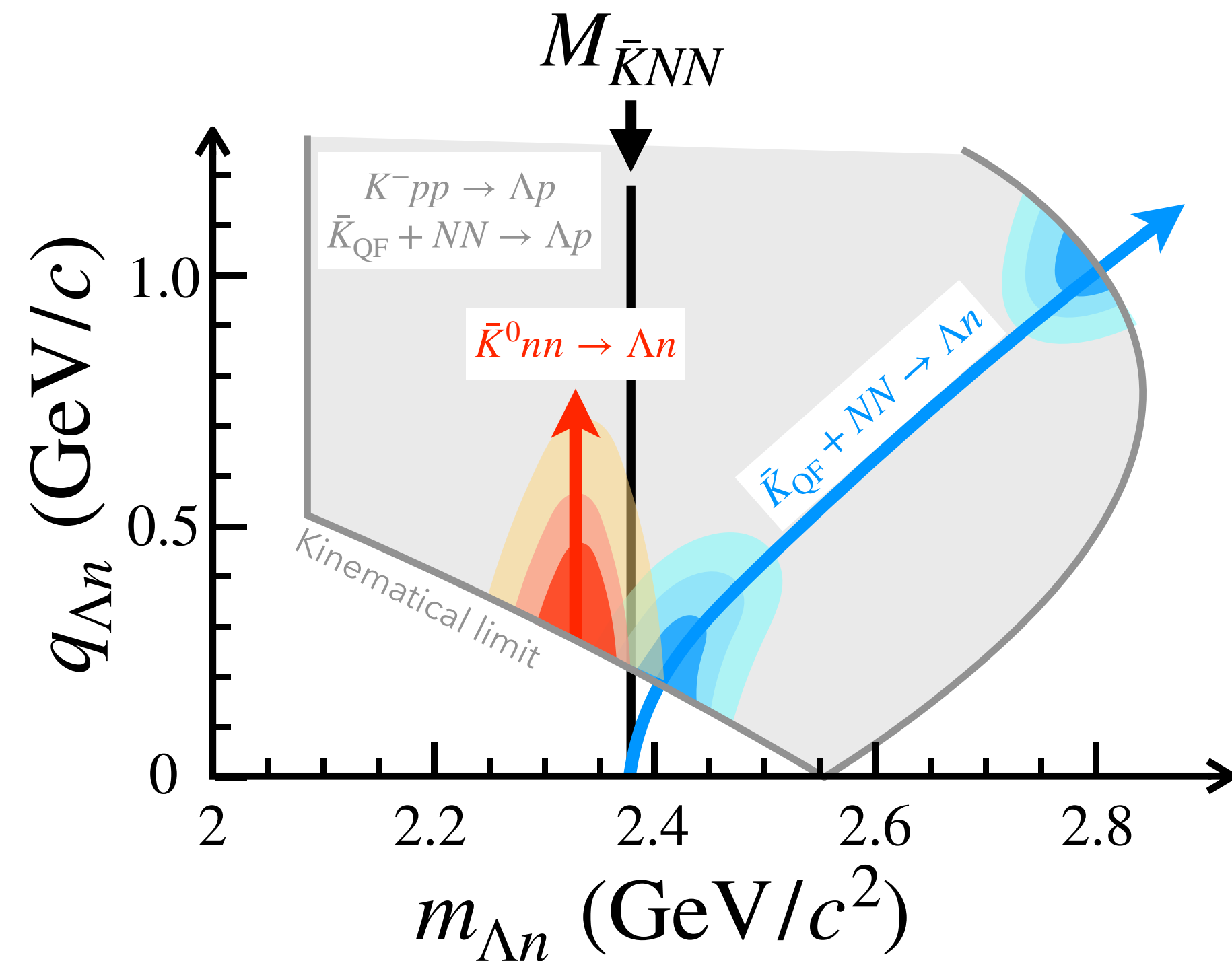
Measurement for Λp pair

*What we did in E15



K^-pp can be seen.

Measurement for Λn pair

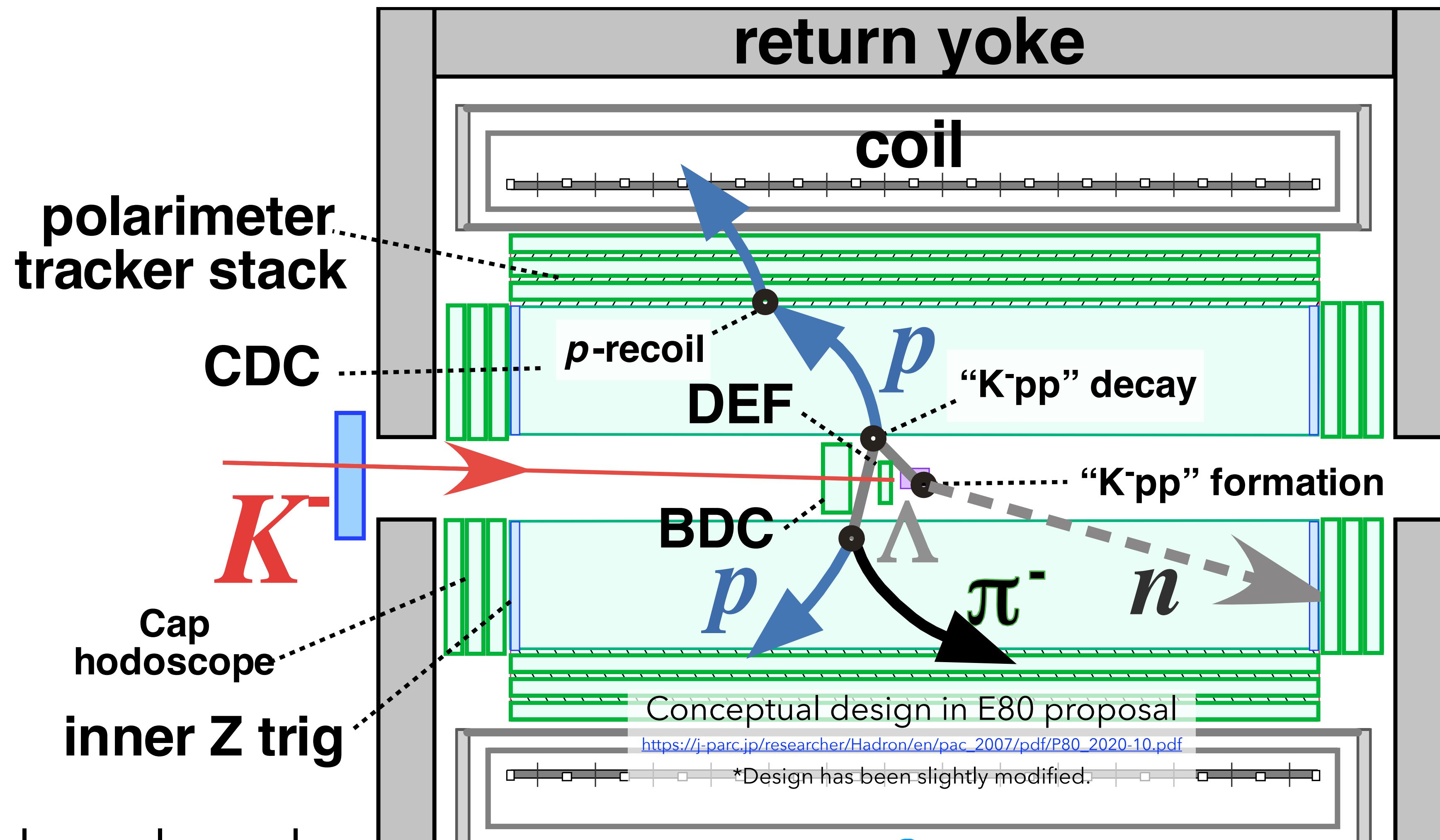


$\bar{K}^0 nn$ can be seen.

K^-pp & $\bar{K}^0 nn$ make event concentration in Λp & Λn pair

Experimental setup

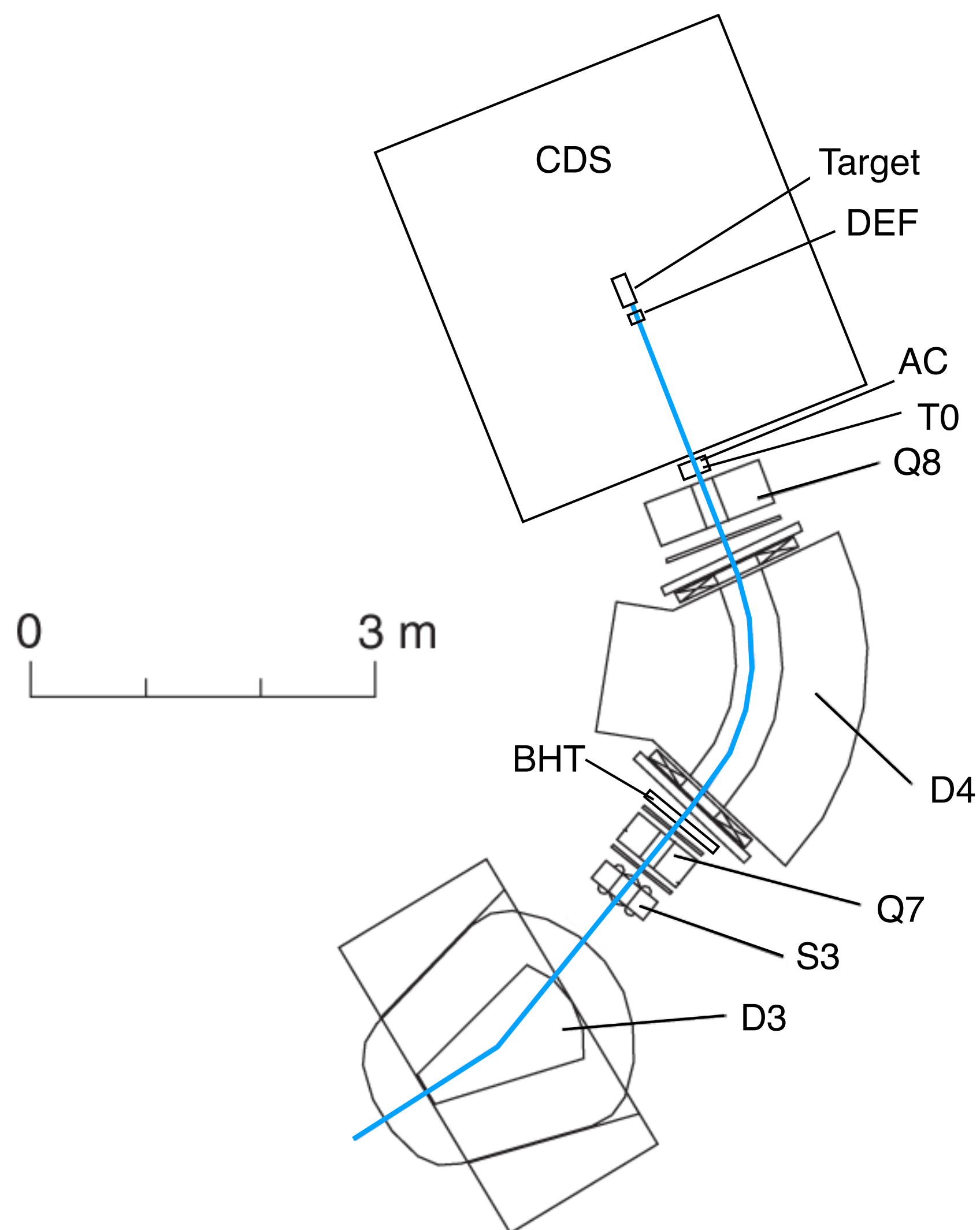
Cylindrical detector system



$\sim 4\pi$ acceptance for charged & neutron

Polarimeter system for p -spin measurement

K1.8BR beam-line



Conceptual design in E80 proposal

https://j-parc.jp/researcher/Hadron/en/pac_2007/pdf/P80_2020-10.pdf

Expected K^- -beam intensity (1 GeV/c)

$\sim 320 \text{ k /spill on target (@ 90kW)}$

Purity
 $K/\pi \sim 1/2$

Expected trigger rate

$\sim 4000 \text{ /spill}$

Rate capability of existing DAQ
>90% efficiency up to $\sim 10000/\text{spill}$

Expected $\bar{K}NN$ yield (reconstructed)

$K^-pp \rightarrow \Lambda p$ ($9.3 \mu\text{b}$) $\bar{K}^0nn \rightarrow \Lambda n$ ($1.2 \mu\text{b}$)

$\sim 6000 \text{ /week}$ $\sim 250 \text{ /week}$

of K^-pp reconstructed in E15
 $\sim 1700 \text{ /(4 weeks)}$

Sensitivity estimation

Sensitivity for $\phi_{\Lambda p}$ -asymmetry

– Formalism of asymmetry reduction factor –

Spin-spin correlation on ϕ -asymmetry

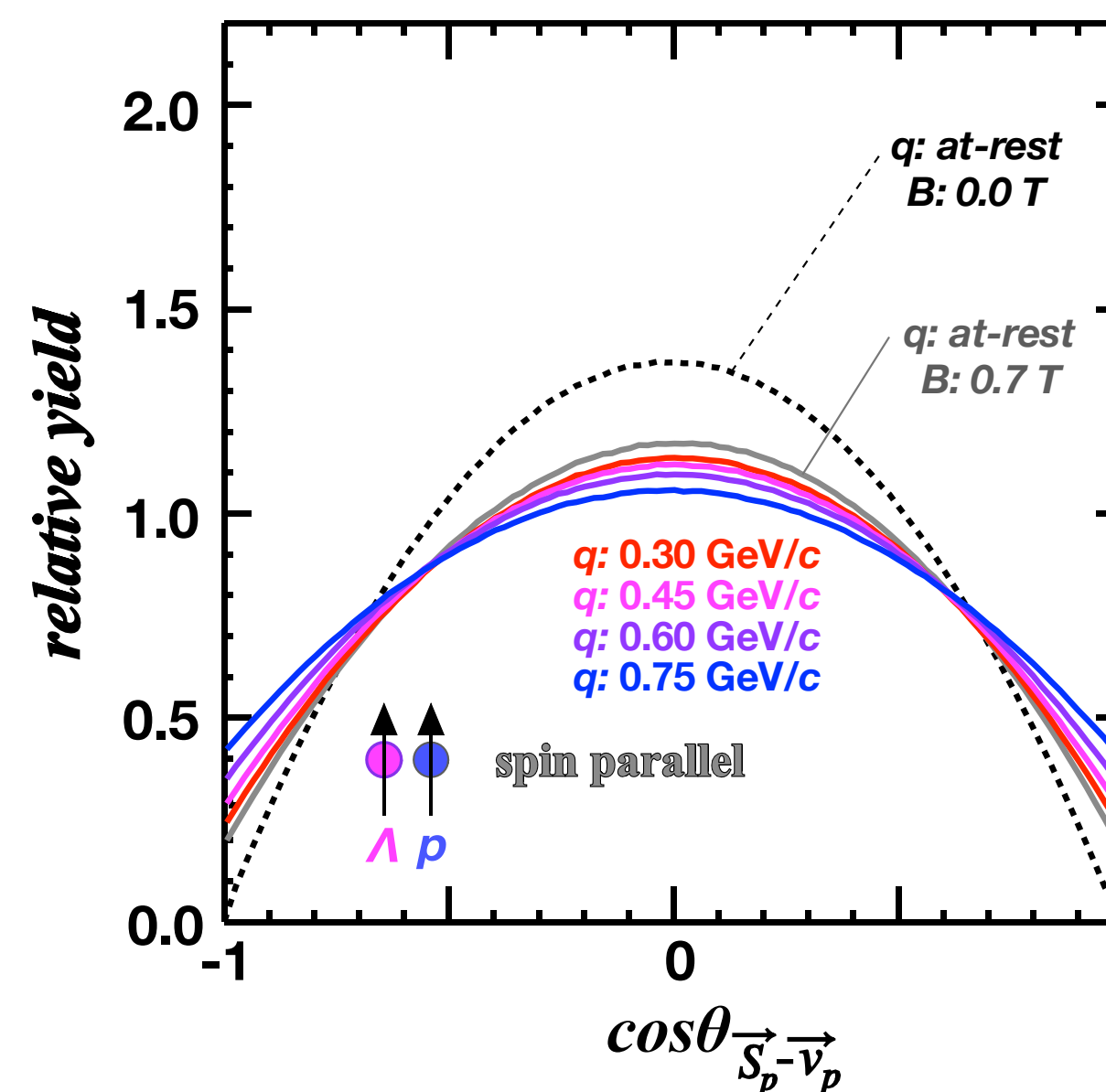
$$N(\phi_{\Lambda p}) = N_0 \cdot (1 + r^{(J^P)} \cdot \alpha_{\Lambda p} \cos \phi_{\Lambda p})$$

$$r^{(J^P)} = \frac{\pi}{4} \alpha_- A_{pC} \int d\Omega f_{\vec{S}_\Lambda} \sin^2 \theta_\Lambda$$

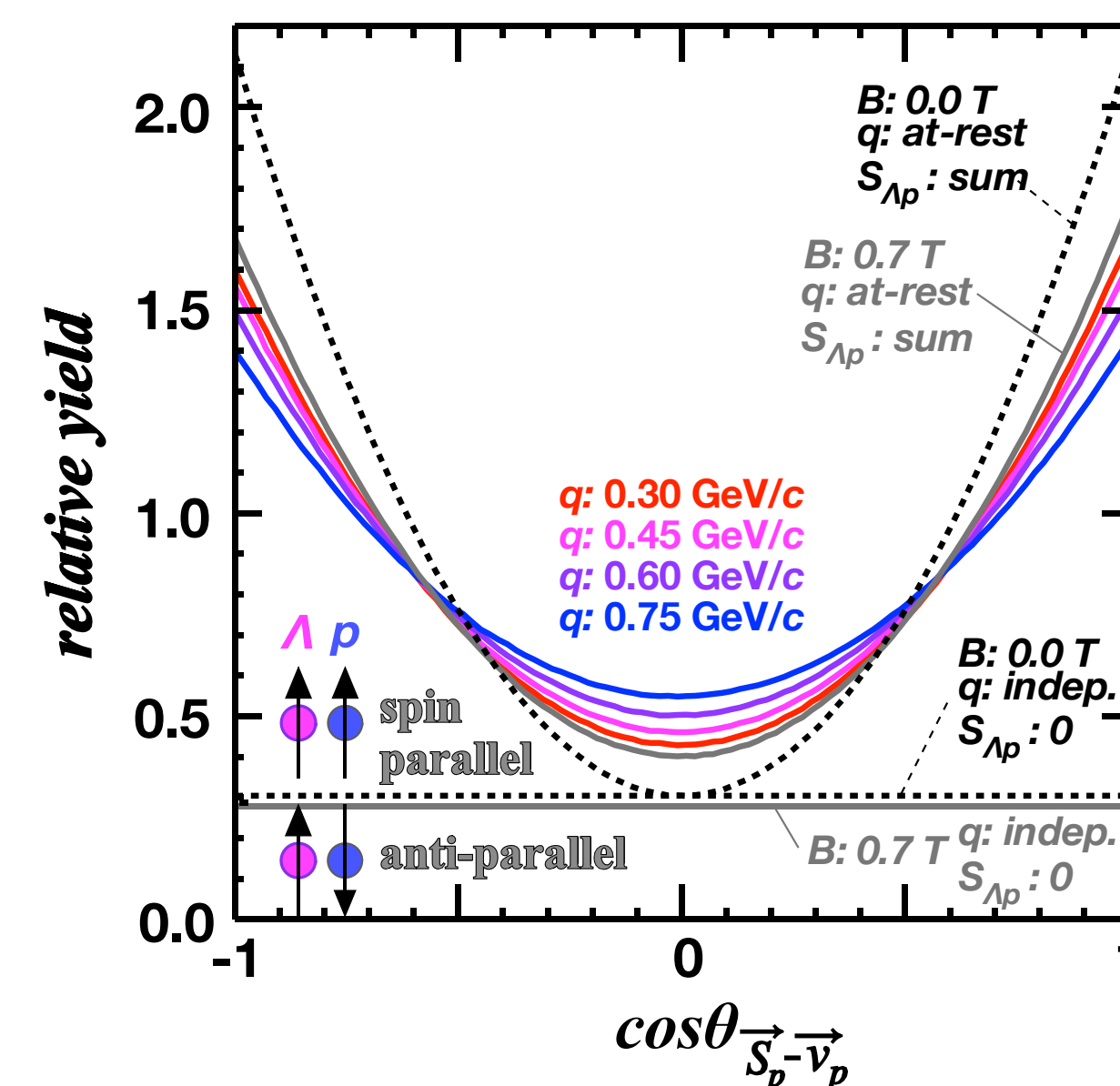
$\alpha_- = 0.72$: Λ asymmetry parameter
 $\langle A_{pC} \rangle \sim 0.3$: Analyzing power

Spin distribution $f_{\vec{S}_\Lambda}$

$J^P = 0^-$



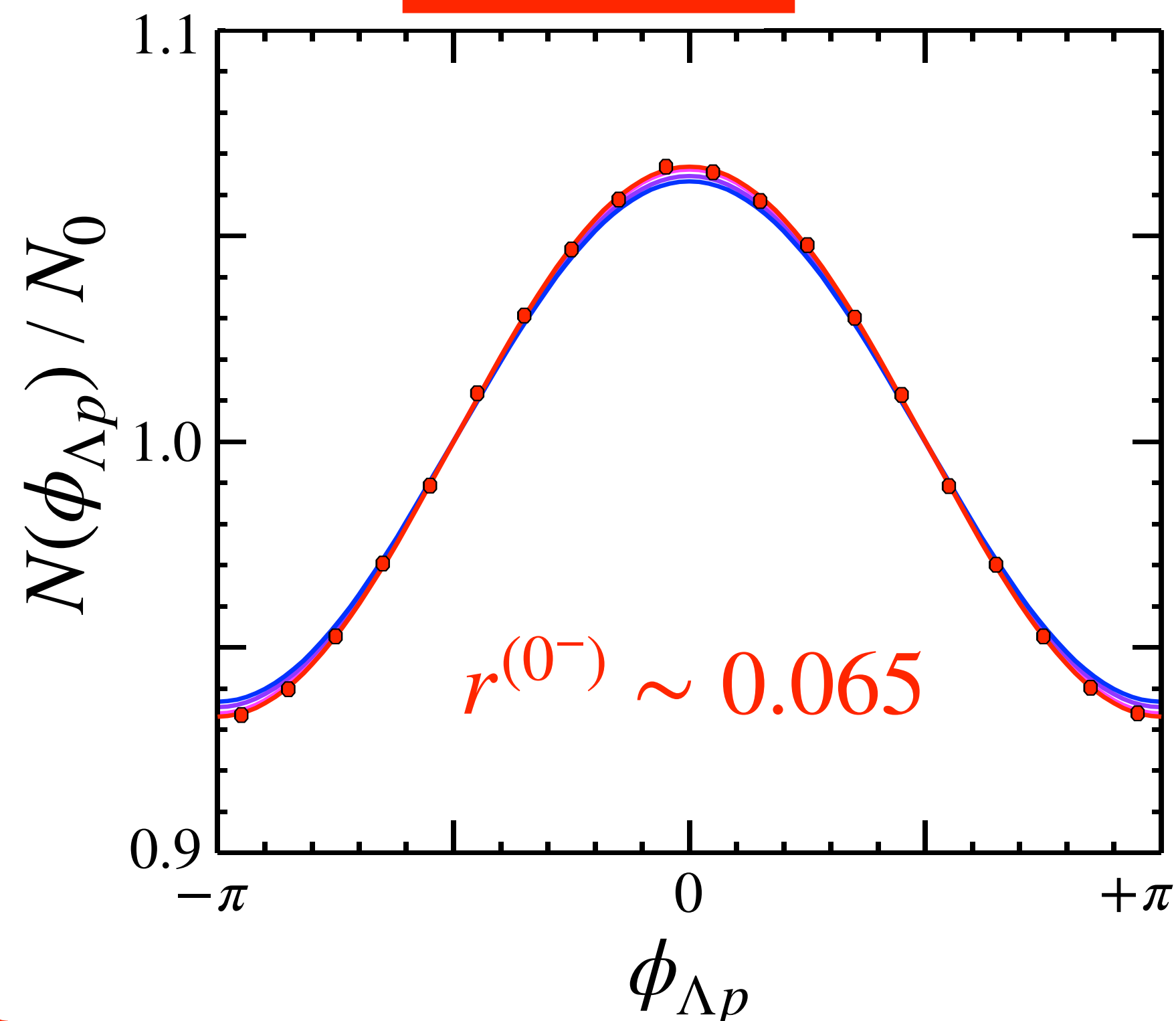
$J^P = 1^-$



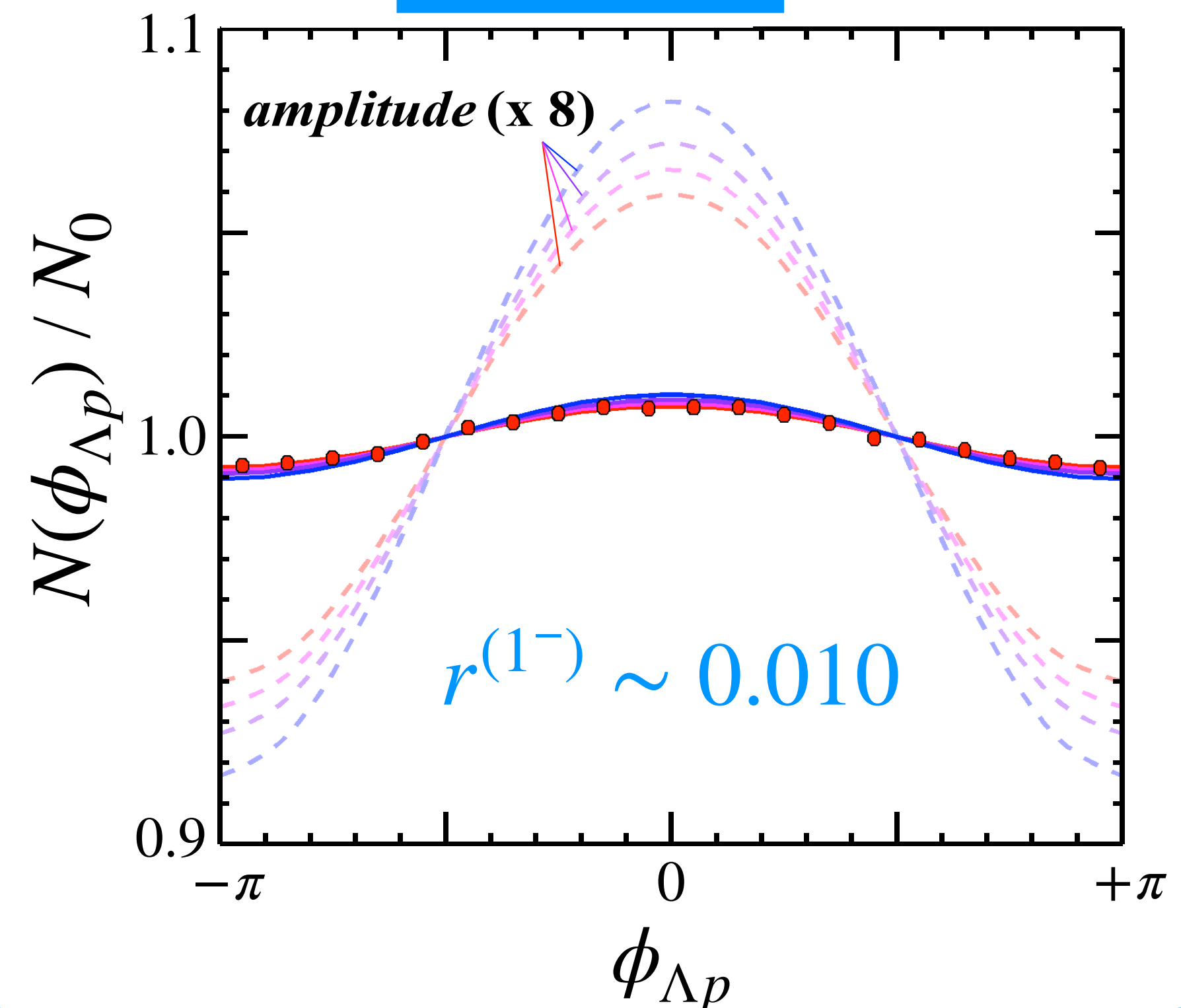
Sensitivity for $\phi_{\Lambda p}$ -asymmetry

– Estimation of observable $\phi_{\Lambda p}$ -asymmetry –

$$J^P = 0^-$$



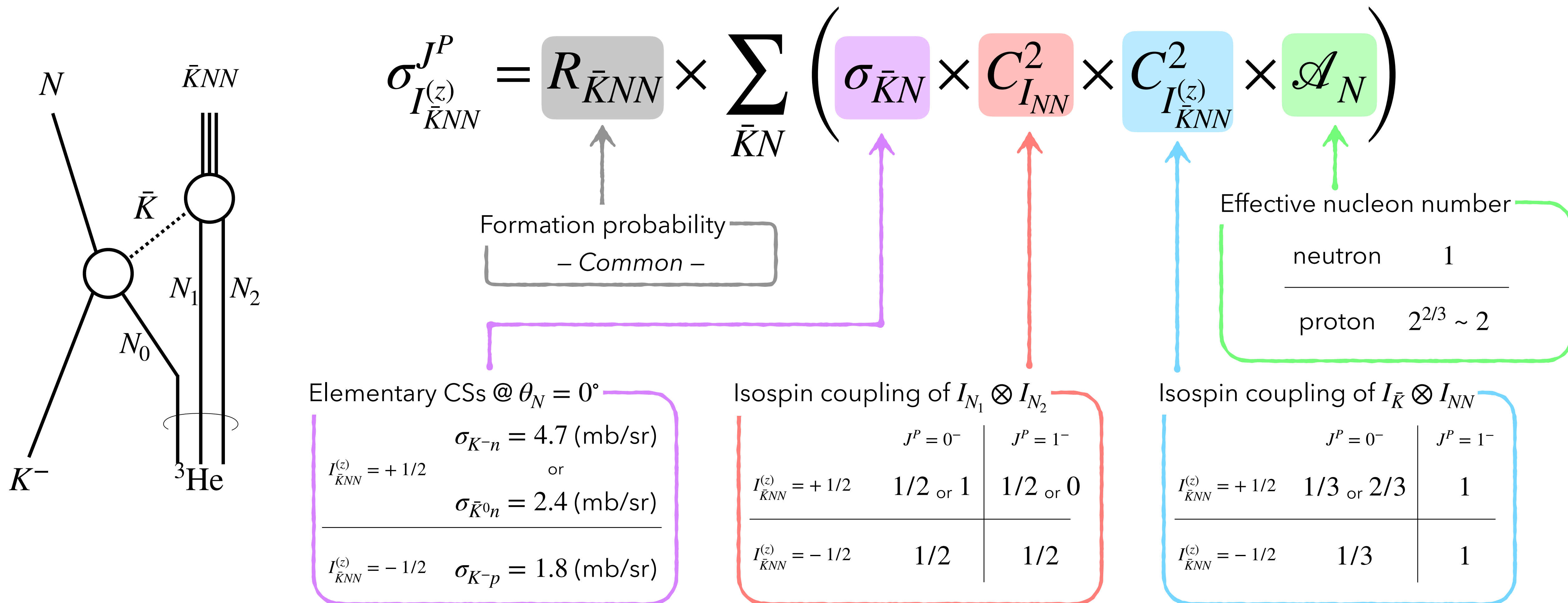
$$J^P = 1^-$$



$r^{(J^P)}$ is large enough to deduce J^P .

Sensitivity for $\bar{K}^0 nn$ search

– Expected formalism of production cross section of $\bar{K}NN$ –



Sensitivity for $\bar{K}^0 nn$ search

– Estimated $\bar{K}^0 nn$ production cross section–

$$\sigma_{K^- pp} \cdot Br_{\Lambda p}$$

$$9.3 \mu\text{b}$$

Measured in E15

$$\sigma_{\bar{K}^0 nn} \cdot Br_{\Lambda n}$$

*Assuming $\mathcal{A}_p = 2^{2/3}$, $Br_{\Lambda n} = Br_{\Lambda p}$

$$J^P = 0^-$$

$$1.2 \mu\text{b}$$

$$J^P = 1^-$$

$$7.0 \mu\text{b}$$

$\sigma_{\bar{K}^0 nn}$ is large enough to see a signal.

$\sigma_{\bar{K}^0 nn} / \sigma_{K^- pp}$ is also a good indicator for J^P .

Beam-time request

8 weeks $\otimes$ 90 kW

(648 kW·week @ 90% up-time)

Ready to start in 2025

Additional tasks from E80

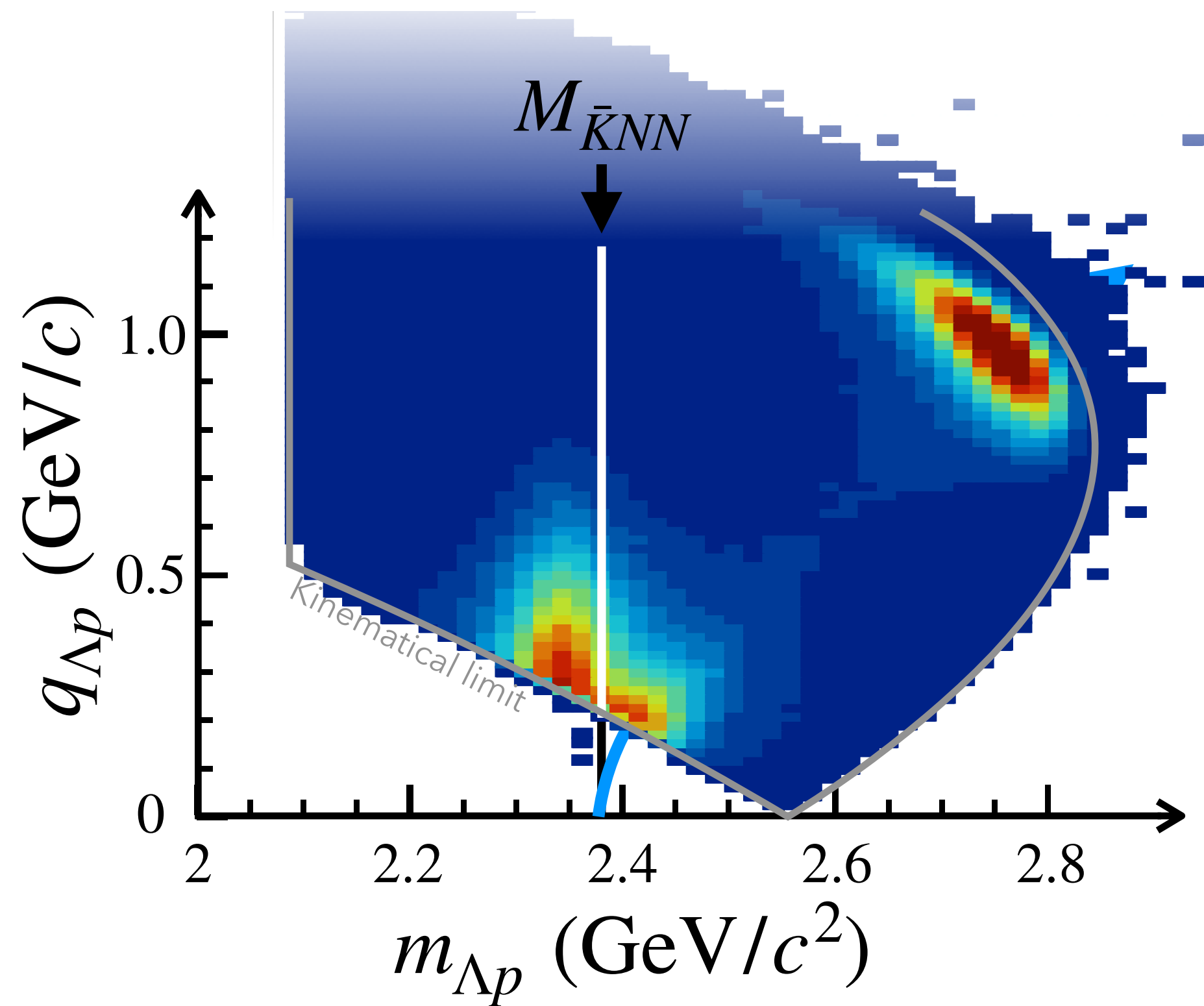
To construct polarimeter system

To purchase additional 300L gaseous ^3He

Expected spectra

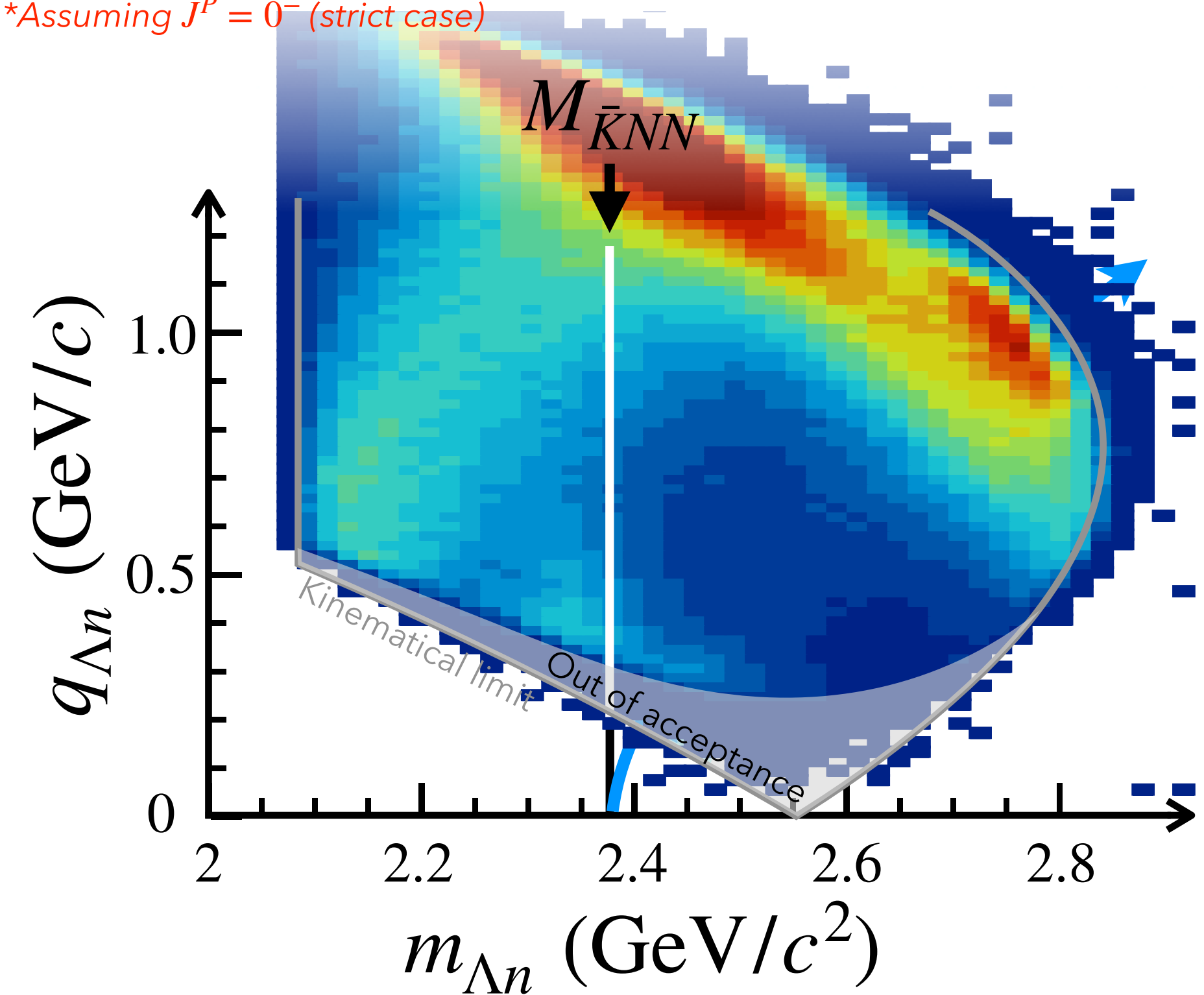
– To detect K^-pp & $\bar{K}^0nn$ –

Measurement for Λp pair



Measurement for Λn pair

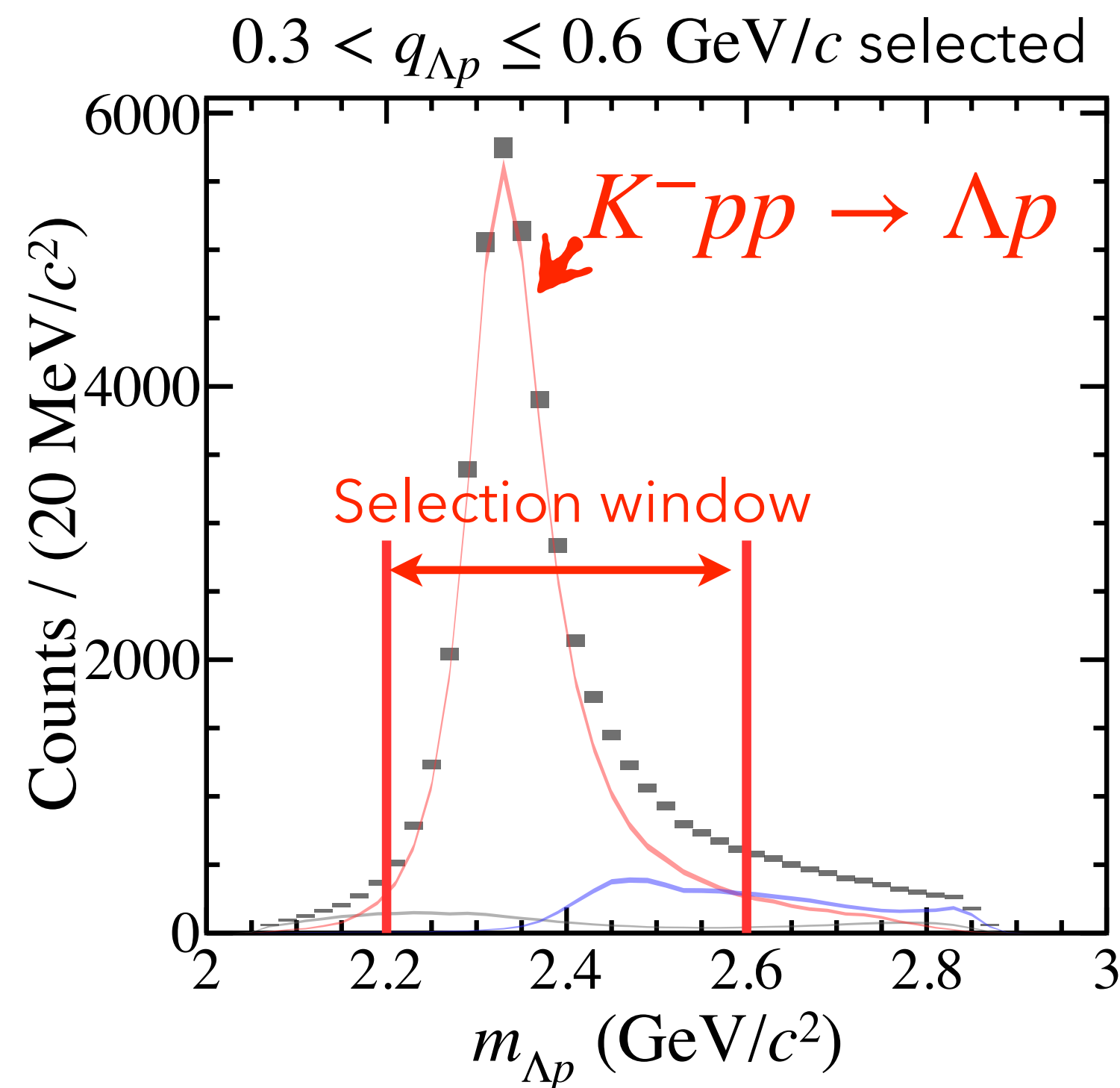
*Assuming $J^P = 0^-$ (strict case)



Expected spectra

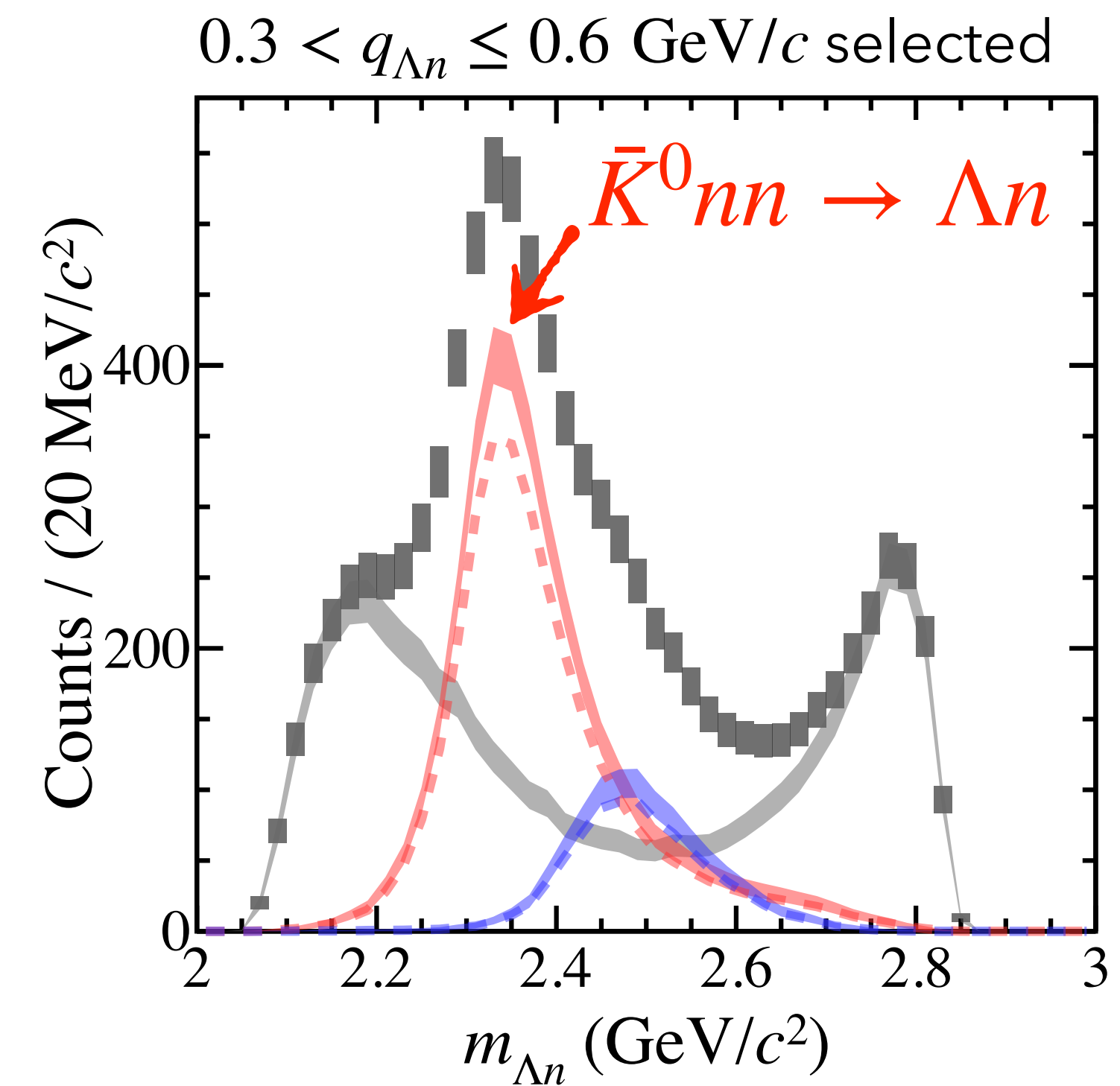
– To detect K^-pp & $\bar{K}^0nn$ –

Measurement for Λp pair



Measurement for Λn pair

*Assuming $J^P = 0^-$ (strict case)

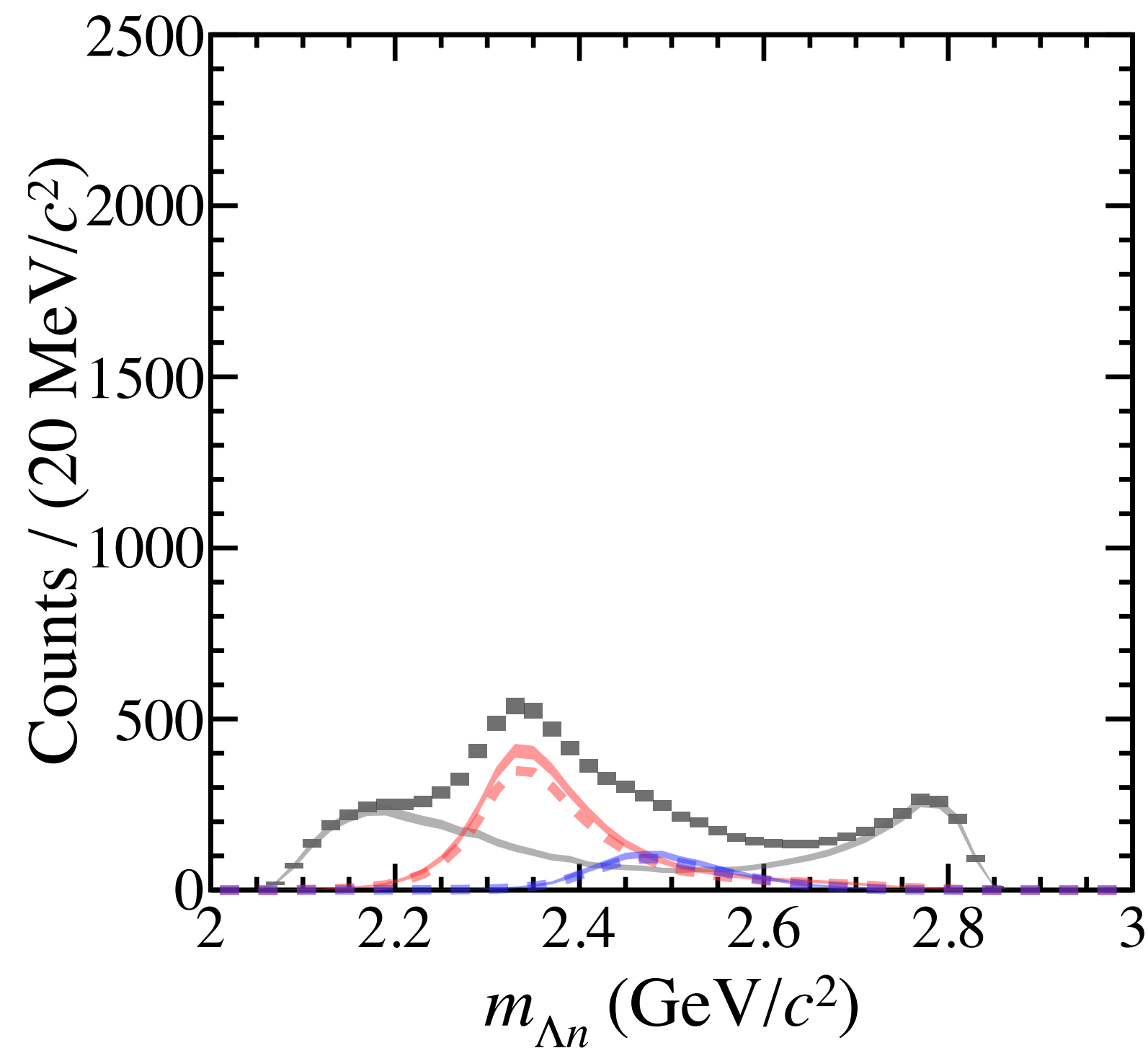


We would observe distinct peaks of both K^-pp & $\bar{K}^0nn$

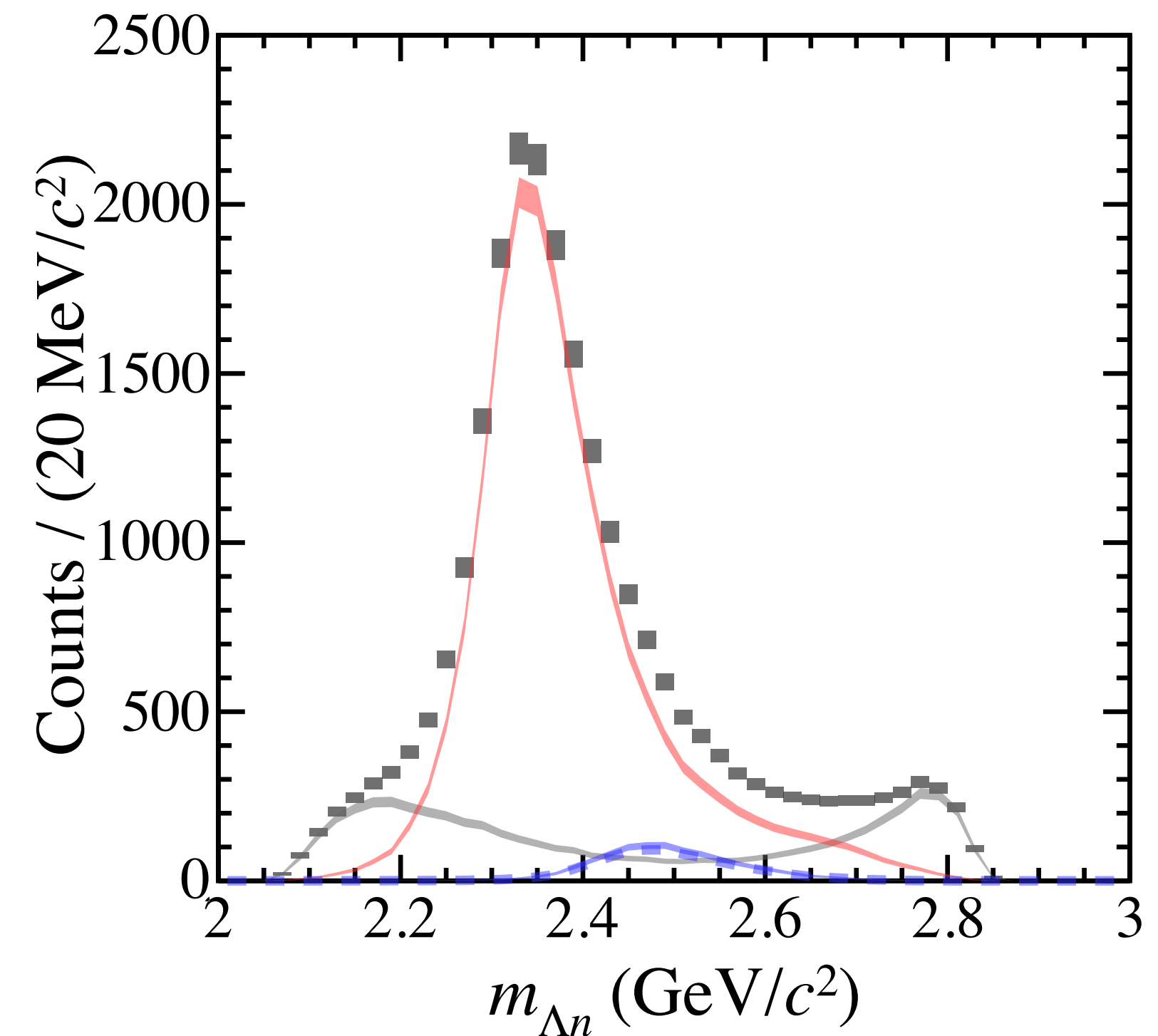
Expected spectra

– To measure $\sigma_{\bar{K}^0 nn}$ for J^P determination–

$J^P = 0^-$



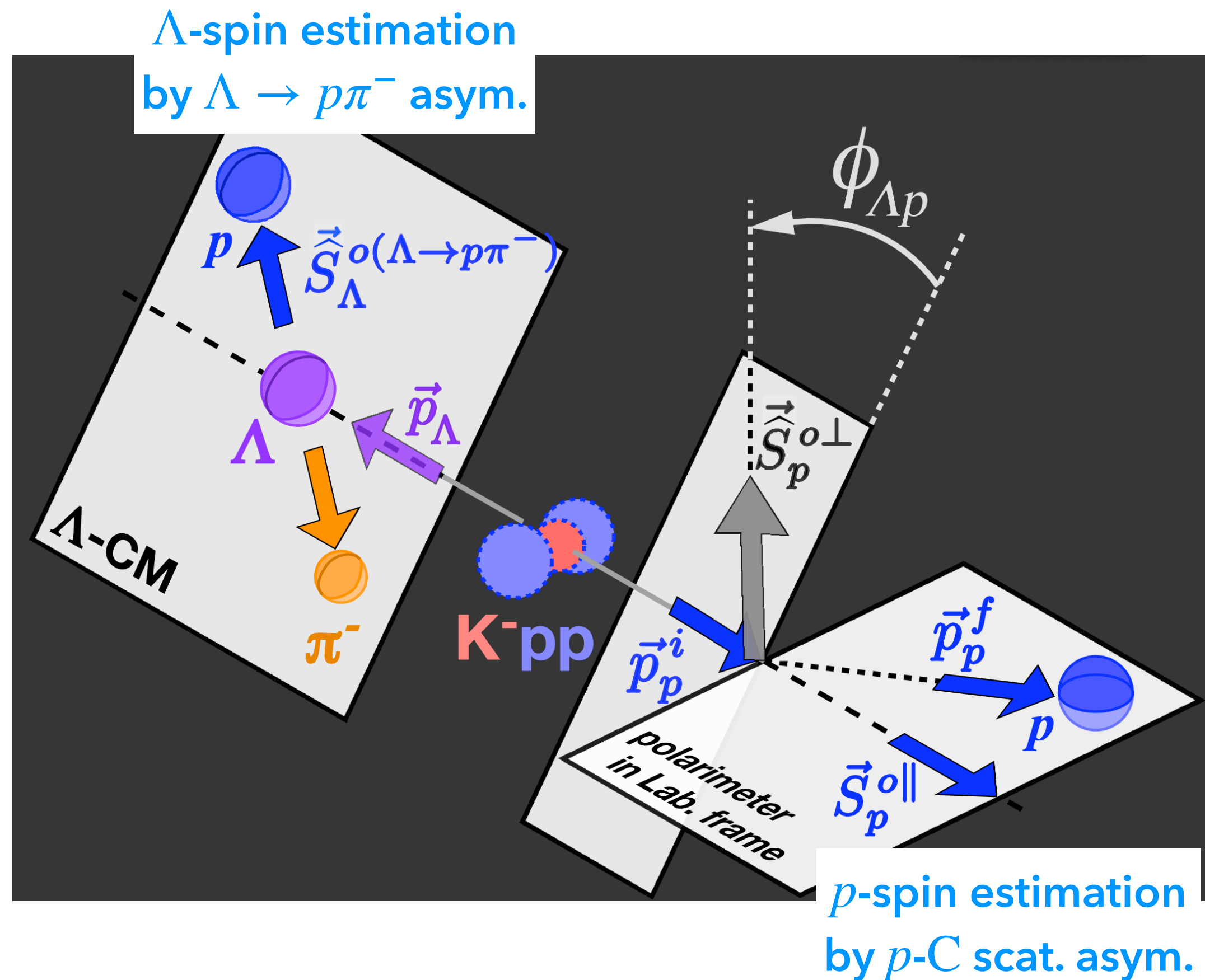
$J^P = 1^-$



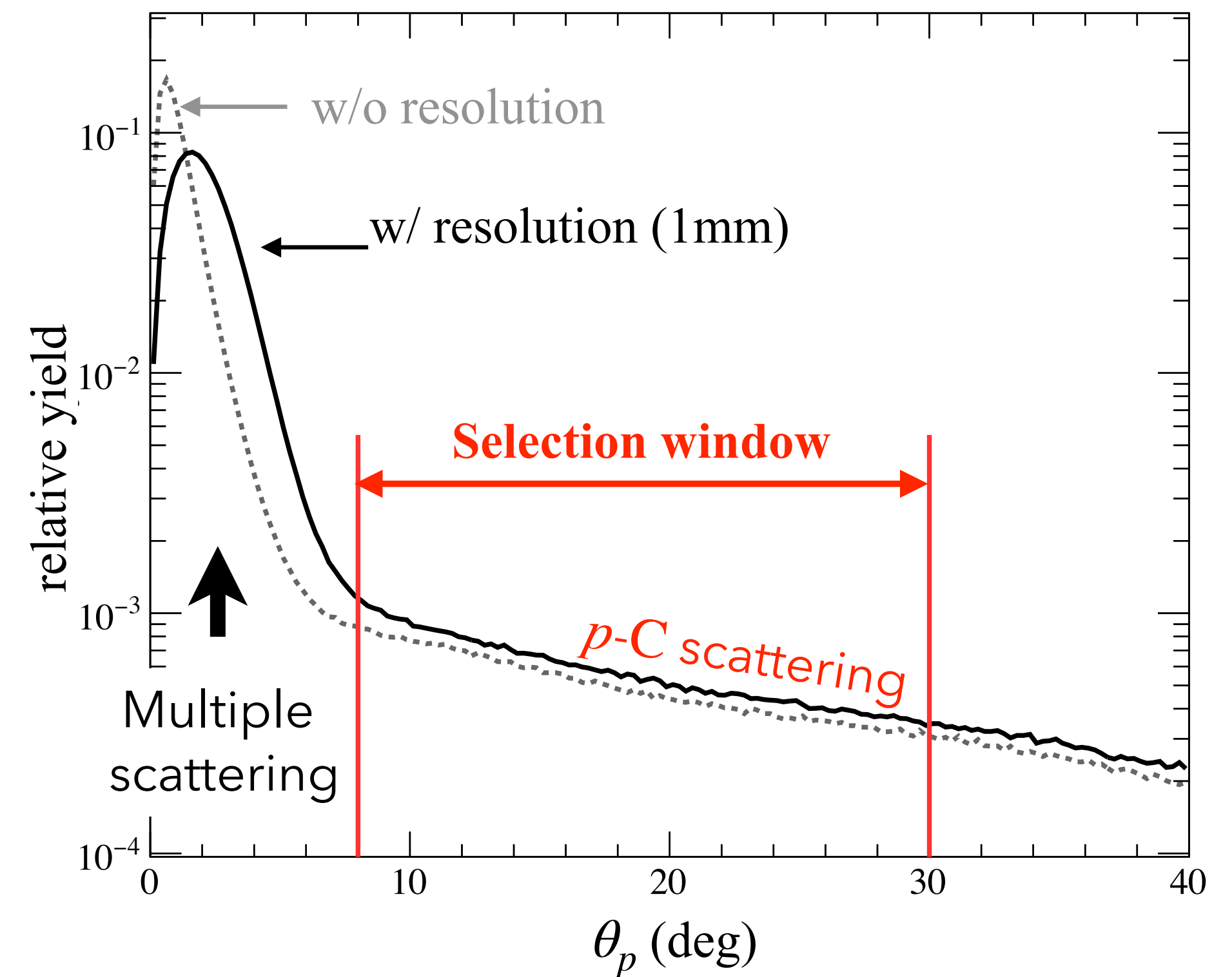
To measure $\sigma_{\bar{K}^0 nn}$ will help us to deduce J^P precisely.

Expected spectra

– To measure $\phi_{\Lambda p}$ -asymmetry for J^P determination –



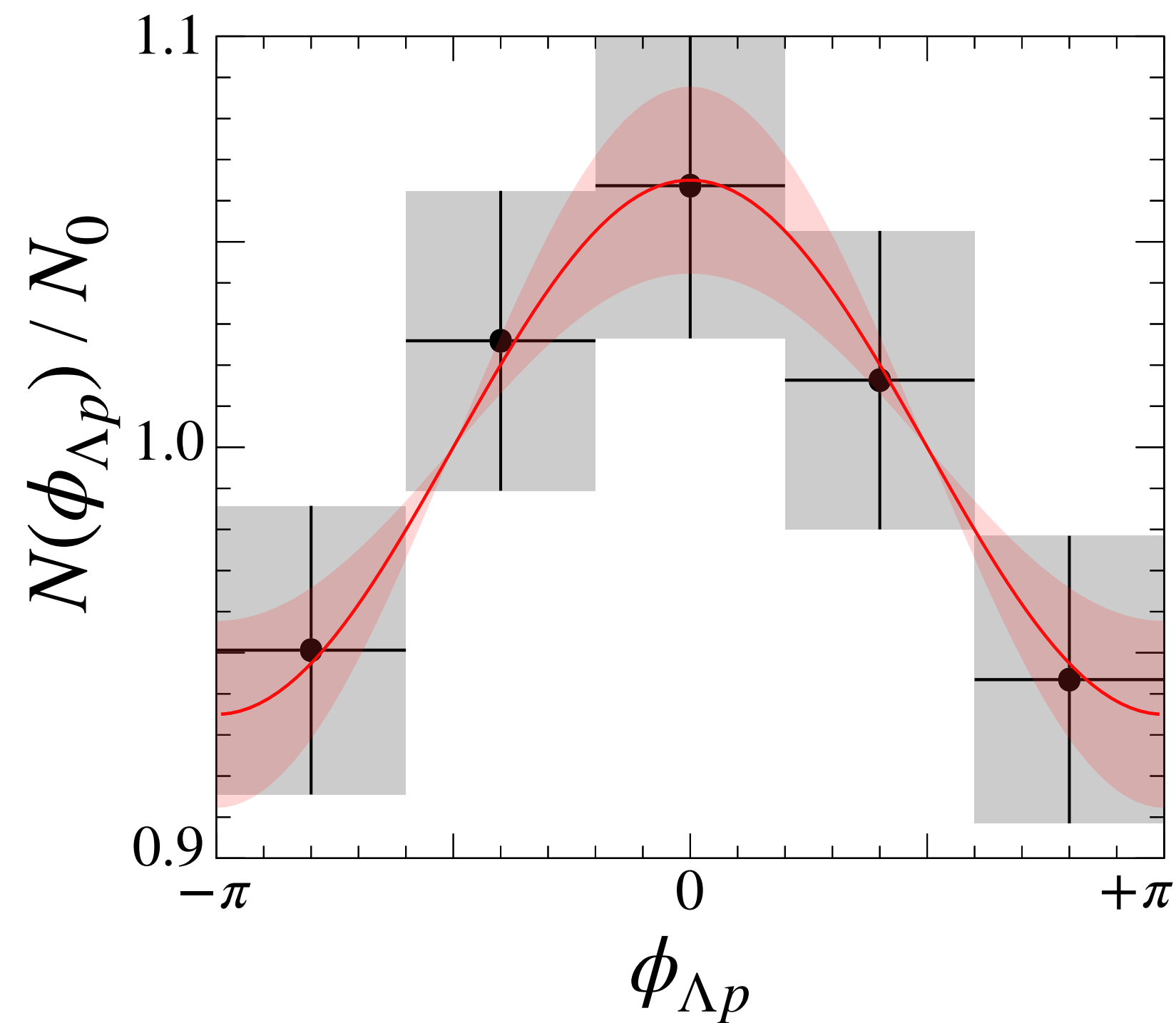
Event selection for p -C scattering



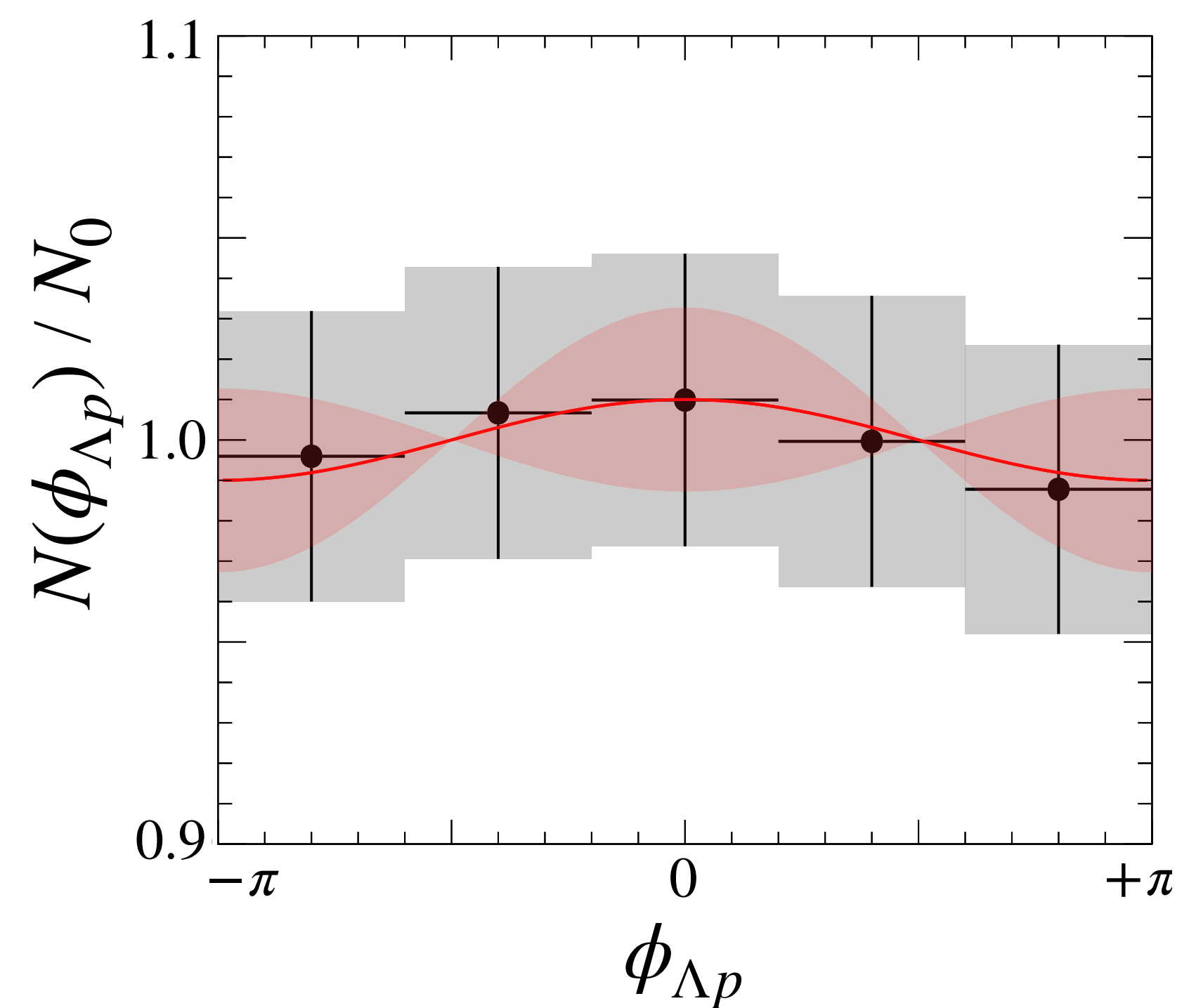
Expected spectra

– To measure $\phi_{\Lambda p}$ -asymmetry for J^P determination –

$J^P = 0^-$



$J^P = 1^-$



We would exclude $J^P = 1^-$ with **95%** confidence level from **only $\phi_{\Lambda p}$ -asym.**

Collaboration

Experimentalists



H. Asano, K. Itahashi, M. Iwasaki, Y. Ma, R. Murayama, H. Ota, F. Sakuma, T. Yamaga



T. Hashimoto, K. Tanida



H. Ohnishi, Y. Sada, C. Yoshida



T. Akaishi



T. Nagae



K. Inoue, S. Kawasaki, H. Noumi, K. Shiotori



M. Bazzi, A. Clozza, C. Curceanu, C. Guaraldo, M. Iliescu, M. Miliucci, A. Scordo, D. Sirghi, F. Sirghi



H. Fujioka



M. Iio, S. Ishimoto, K. Ozawa, S. Suzuki



J. Marton, H. Shi, M. Tuechler, E. Widmann, J. Zmeskal

Theorists



D. Jido



T. Sekihara

We investigate $\bar{K}NN$ both **experimentally** & **theoretically**.
Especially, to deduce the spacial size from q -dependence

Summary

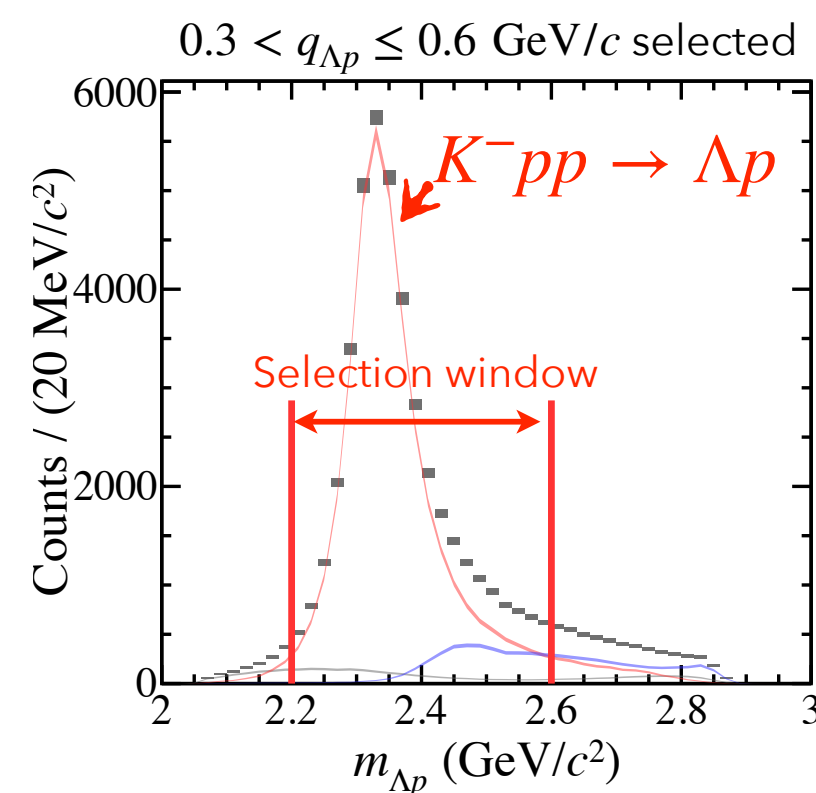
– For stage-1 approval –

Expected results

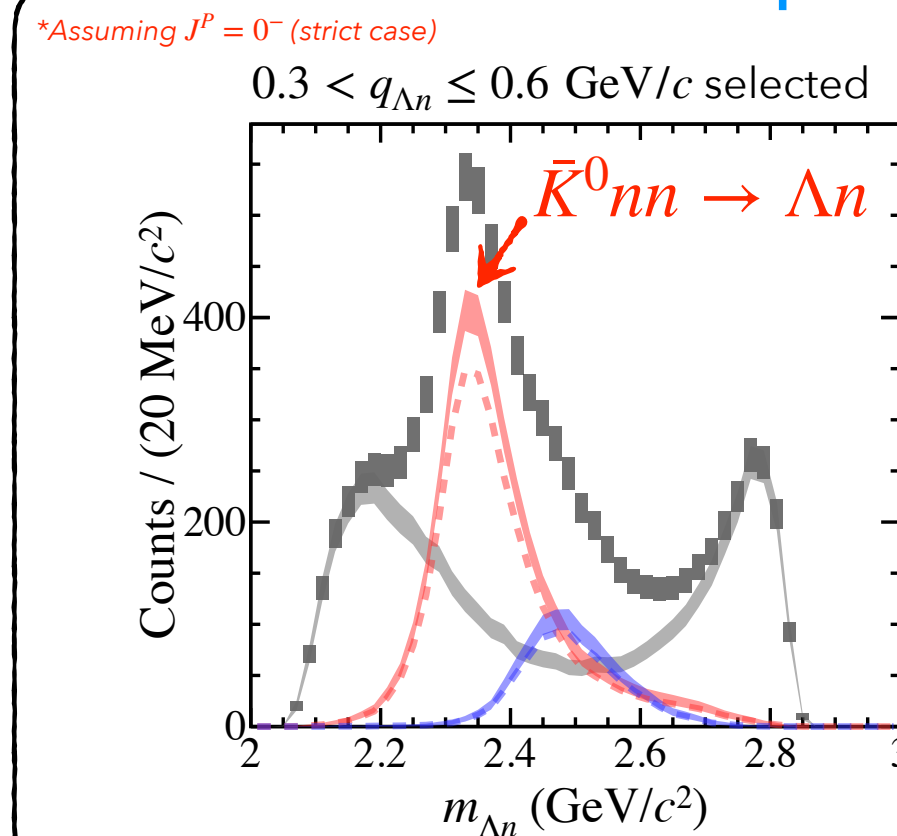
8 weeks $\otimes$ 90 kW @ K1.8BR

To confirm $I = 1/2 \bar{K}NN$ doublet

Measurement for Λp pair

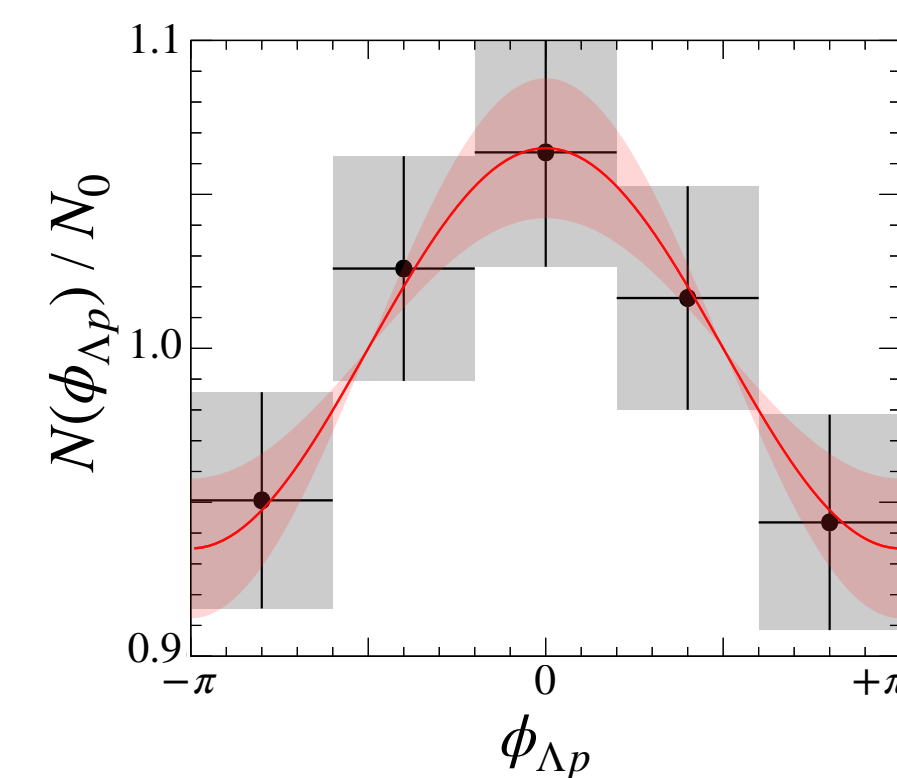


Measurement for Λn pair

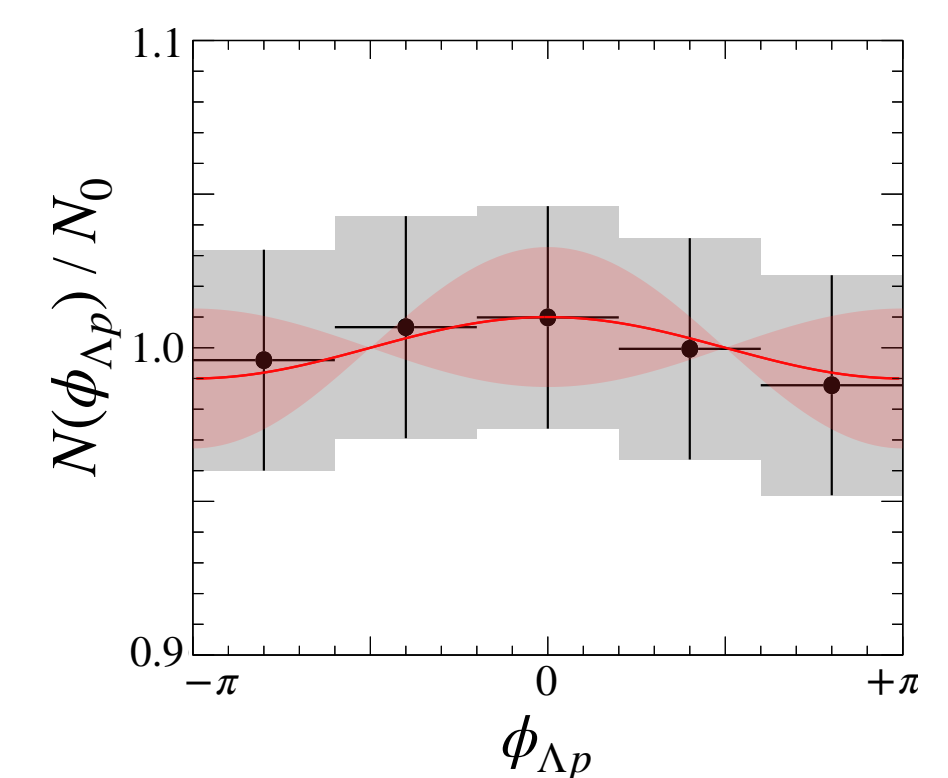


To determine J^P of $\bar{K}NN$

$J^P = 0^-$



$J^P = 1^-$

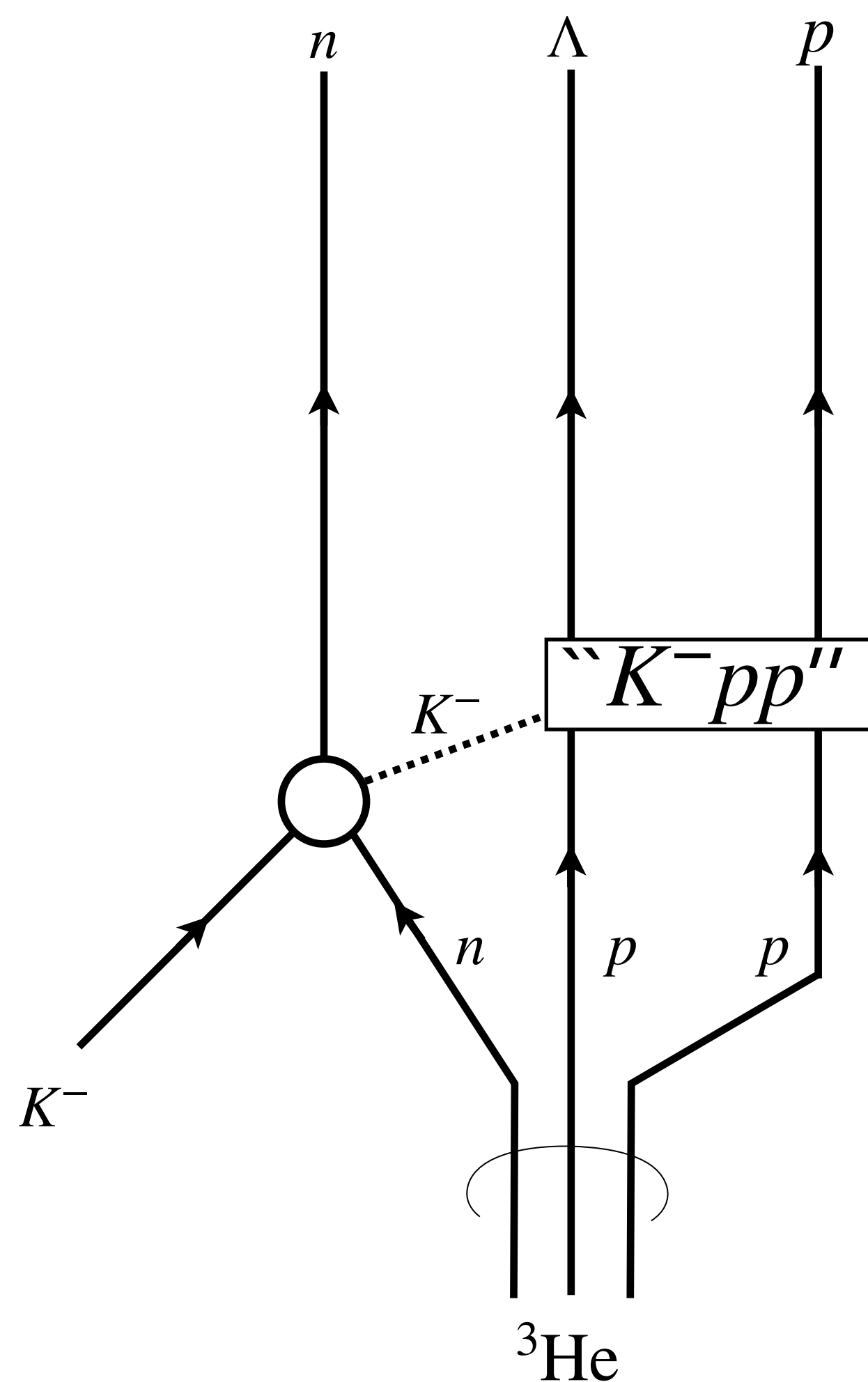
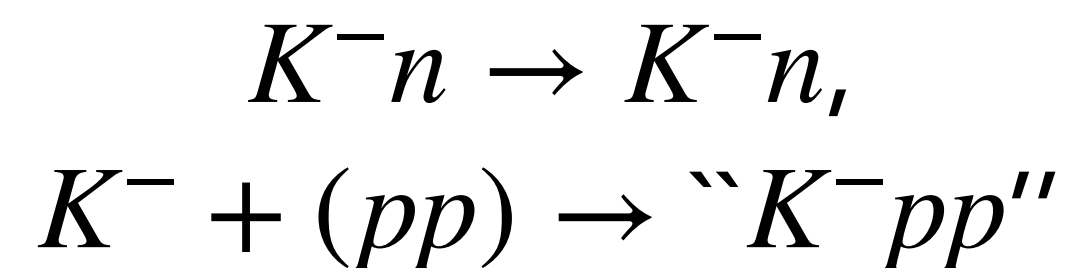


We wish to open a new doorway to access high-density matter

Thank you for your attention!

Backup

The diagram illustrates the addition of two $1/2 \times 1/2$ tensors (represented by red boxes) to form a $1 \times 1/2$ tensor (represented by a gray box). The first $1/2 \times 1/2$ tensor has elements 1 and $+1$. The second $1/2 \times 1/2$ tensor has elements $+1/2 + 1/2$ and 1 . The resulting $1 \times 1/2$ tensor has elements $3/2$ and $+3/2$. The diagram also shows the resulting $1 \times 1/2$ tensor being added to another $1 \times 1/2$ tensor (represented by a gray box) to form a $1 \times 1/2$ tensor (represented by a red box). The resulting $1 \times 1/2$ tensor has elements 1 and $+1$. The diagram also shows the resulting $1 \times 1/2$ tensor being added to another $1 \times 1/2$ tensor (represented by a gray box) to form a $1 \times 1/2$ tensor (represented by a red box). The resulting $1 \times 1/2$ tensor has elements 1 and $+1$.



$$I_{NN} = \overset{p}{|1/2, +1/2\rangle} \otimes \overset{p}{|1/2, +1/2\rangle} = |1, +1\rangle$$

; Only $J^P = 0^-$ can be produced.

$$I_{\bar{K}NN} = |1, +1\rangle \otimes |1/2, -1/2\rangle$$

$$= \sqrt{\frac{1}{3}} |3/2, +1/2\rangle_{I_{NN}=1} + \sqrt{\frac{2}{3}} |1/2, +1/2\rangle_{I_{NN}=1}$$

$$J^P = 0^-$$

$J^P = 0^-$	$J^P = 1^-$
$1 \times \frac{2}{3}$	0

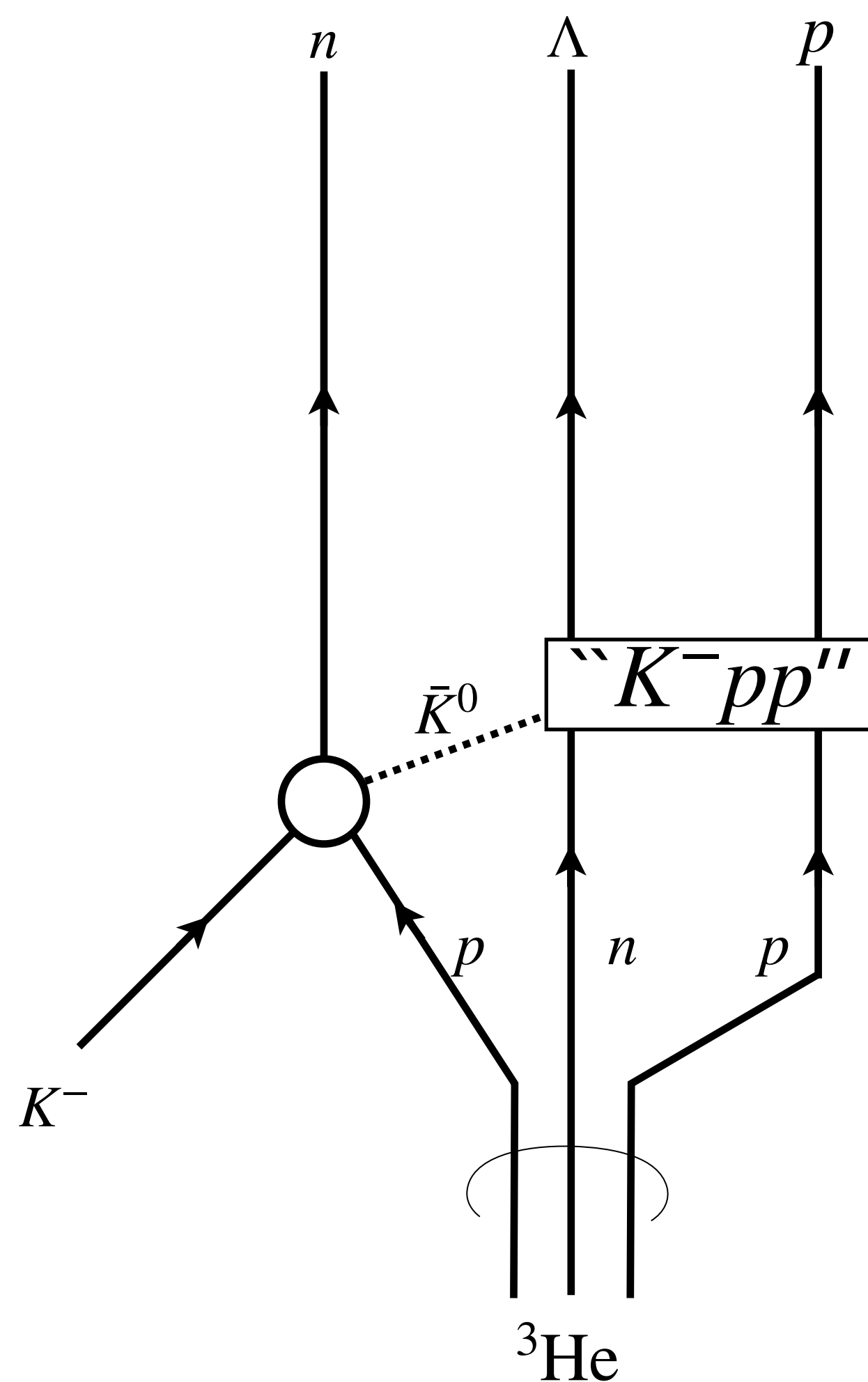
Production of `` K^-pp ''

$1/2 \times 1/2$		1		
	+1	1	0	
+1/2+1/2	1	0	0	
	+1/2 -1/2	1/2	1/2	1
-1/2 +1/2	1/2 -1/2	-1		
	-1/2 -1/2	1		

$1 \times 1/2$		3/2		
	+3/2	3/2	1/2	
+1 +1/2	1	1/2 +1/2		
	+1 -1/2	1/3	2/3	
0 +1/2	2/3 -1/3			

$$K^-p \rightarrow \bar{K}^0 n,$$

$$\bar{K}^0 + (pn) \rightarrow \text{``}K^-pp\text{''}$$



$$I_{NN} = |1/2, +1/2\rangle \otimes |1/2, -1/2\rangle = \sqrt{\frac{1}{2}} (|1,0\rangle + |0,0\rangle)$$

$$I_{\bar{K}NN} = \sqrt{\frac{1}{2}} (|1,0\rangle + |0,0\rangle) \otimes |1/2, +1/2\rangle$$

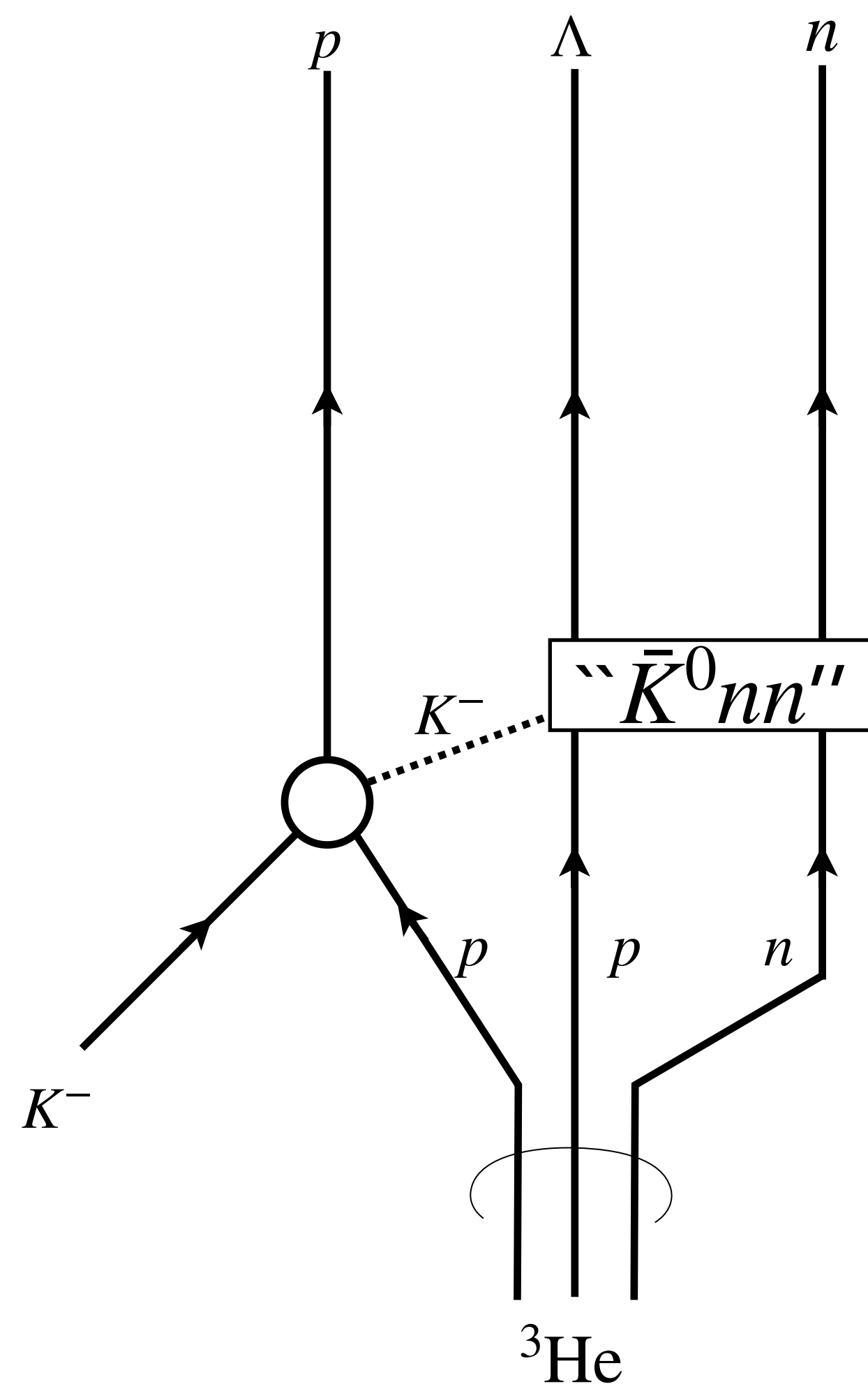
$$= \sqrt{\frac{1}{2}} \left(\sqrt{\frac{2}{3}} |3/2, +1/2\rangle_{I_{NN}=1} - \sqrt{\frac{1}{3}} |1/2, +1/2\rangle_{I_{NN}=1} \right)$$

$$+ \sqrt{\frac{1}{2}} |1/2, +1/2\rangle_{I_{NN}=0}$$

$$J^P = 0^-$$

$$J^P = 1^-$$

$J^P = 0^-$	$J^P = 1^-$
$\frac{1}{2} \times \frac{1}{3}$	$\frac{1}{2}$



$$I_{NN} = \overset{p}{|1/2, +1/2\rangle} \otimes \overset{n}{|1/2, -1/2\rangle} = \sqrt{\frac{1}{2}} (|1,0\rangle + |0,0\rangle)$$

$$\begin{aligned}
I_{\bar{K}NN} &= \sqrt{\frac{1}{2}} (|1,0\rangle + |0,0\rangle) \otimes |1/2, -1/2\rangle \\
&= \sqrt{\frac{1}{2}} \left(\sqrt{\frac{2}{3}} |3/2, -1/2\rangle_{I_{NN}=1} + \sqrt{\frac{1}{3}} |1/2, -1/2\rangle_{I_{NN}=1} \right) \\
&\quad + \sqrt{\frac{1}{2}} |1/2, -1/2\rangle_{I_{NN}=0}
\end{aligned}$$

$J^P = 0^-$	$J^P = 1^-$
1 1	1

$J^P = 0^-$	$J^P = 1^-$
$\frac{1}{2} \times \frac{1}{3}$	$\frac{1}{2}$

Expected spectra

- To detect $\bar{K}^0_{nn}$ & study isospin symmetry for $\bar{K}NN$ –

